# ROBUST NORM EQUIVALENCIES AND OPTIMAL PRECONDITIONERS FOR DIFFUSION PROBLEMS 

MICHAEL GRIEBEL, KARL SCHERER, AND MARC ALEXANDER SCHWEITZER


#### Abstract

Additive multilevel methods offer an efficient way for the fast solution of large sparse linear systems which arise from a finite element discretization of an elliptic boundary value problem. These solution methods are based on norm equivalencies for the associated bilinear form using a suitable subspace decomposition. From these, efficient preconditioners for the stiffness matrix can be derived. For the robustness of the resulting iteration schemes, it is crucial that the constants in the norm equivalence do not depend or depend only weakly on the ellipticity constants of the problem. Here we address this question for the model problem $-\nabla \omega \nabla u=f$ with a scalar weight $\omega$.

We prove an upper bound completely independent of the weight $\omega$, whereas our lower bound involves some information about the local variation the coefficient function. The condition on the coefficient is related to the $A_{2}$ and $A_{\infty}$ Muckenhoupt classes from approximation theory. It is sufficient to require $\omega \in A_{2} \subset A_{\infty}$, however, the results even hold for a slightly larger class of functions.


## 1. Introduction

The solution of large sparse linear systems arising from the discretization of a partial differential equation (PDE) is an essential ingredient in many scientific computations. The ever growing demand for efficient solvers led to the development of multigrid methods in the 1970s [4, 5, 13, 14] and multilevel preconditioning techniques in the 1980s [22]. Much research work was devoted to the question of optimal complexity; i.e., to show that the number of operations necessary to obtain the solution up to a prescribed accuracy is proportional to the number of unknowns of the linear system. However, the convergence behavior of these classical schemes is still strongly dependent on the coefficients of the considered PDE. This is the so-called robustness problem of multilevel solvers. It is one major reason which somewhat limits the applicability of classical multigrid methods and multilevel preconditioners to real world problems. Several extensions of these methods e.g. via the use of more complicated smoothing schemes or through the use of operator-dependent or matrix-dependent transfer operators $[1,10,24,25]$ in the so-called black-box multigrid method have been proposed over the years to overcome the robustness problem. Currently the most successful approach is the algebraic multigrid (AMG) method $[6,7,8,9,12,16,19]$ which further generalizes the black-box multigrid idea.

[^0]However, the design and implementation of an AMG method is rather involved and there is no satisfying theoretical foundation which proves the robustness of AMG. Hence, the development of a provably robust multilevel solver or preconditioner for general second order PDEs remains an open problem even today.

In this paper we focus on the robustness issue for diffusion problems in two space dimensions involving a scalar diffusion coefficient $\omega: \mathbb{R}^{2} \rightarrow \mathbb{R}$; i.e., we consider the model problem

$$
-\nabla \omega \nabla u=f \text { in } \Omega \subset \mathbb{R}^{2}
$$

We discretize our model problem using linear finite elements on a sequence of uniformly refined triangulations. Furthermore, we only require the weight $\omega$ to be in the Muckenhoupt class $A_{2} \subset A_{\infty}$; in fact, the results hold even for a slightly larger class of functions.

The main ingredients of our analysis are the use of certain weighted projection operators and auxiliary level-dependent elliptic projection operators. The tools employed in the proof are the Hardy-inequality, a hybrid Bernstein-type inequality and a generalized duality technique. We limit ourselves to the two-dimensional case using linear finite elements only, however, we believe that the results can be extended to three space dimensions and multi-linear elements.

In this paper we establish the equivalence

$$
a(u, u) \leq \tilde{C} \sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2} \leq \hat{C} a(u, u)
$$

for a certain weighted norm where the constants $\tilde{C}$ and $\hat{C}$ depend on the initial triangulation and in general involve some information about the local variation of the coefficient function $\omega$. For the standard test case of a piecewise constant coefficient function with maximal jump of height $\epsilon^{-1}$, we obtain an optimal norm equivalence with $\hat{C} \approx \ln \left(\epsilon^{-1 / 2}\right)$. Note that the location of the jumps, their number or their frequency is not restricted. Hence, unlike other articles concerned with the development of robust solvers $[11,15,20]$ we do not require that the jumps must be resolved on any particular level.

The remainder of the paper is organized as follows: First, we introduce the notation and the employed projection operators in $\S 2$. Then, we present the main result of the paper in $\S 3$. Here, we establish two robust norm equivalencies for the considered model problem using a linear finite element discretization on a sequence of uniformly refined triangulations. We begin with the derivation of an optimal and robust upper bound for the bilinear form using a hybrid Bernstein-type estimate and a Hardy-inequality. Then, we establish a lower bound for the bilinear form under a weak assumption on the coefficient function $\omega$ using a duality technique and a Hardy-inequality. Finally, we discuss the conditions imposed on the coefficient functions and show that all weights $\omega$ in the Muckenhoupt class $A_{2} \subset A_{\infty}$ fulfill the requirements. In $\S 4$ we present the construction of a preconditioner of BPX-type based on our robust and optimal norm equivalence before we conclude with some remarks in $\S 5$.

## 2. Prerequisites

Let us introduce some notation which we will use throughout this paper. We will consider a sequence of uniformly refined triangulations

$$
\begin{equation*}
\mathcal{T}_{0} \subset \mathcal{T}_{1} \subset \cdots \subset \mathcal{T}_{J} \tag{2.1}
\end{equation*}
$$

and the associated sequence of piecewise linear finite element spaces

$$
\begin{equation*}
\mathcal{V}_{0} \subset \mathcal{V}_{1} \subset \cdots \subset \mathcal{V}_{J} \tag{2.2}
\end{equation*}
$$

Let $\mathcal{N}_{j}$ denote the set of all vertices of the triangulation $\mathcal{T}_{j}$ and define

$$
\begin{equation*}
\mathcal{R}_{0}:=\mathcal{N}_{0}, \quad \mathcal{R}_{j}:=\mathcal{N}_{j} \backslash \mathcal{N}_{j-1}, \quad \text { and } \mathcal{N}=\mathcal{N}_{J} \tag{2.3}
\end{equation*}
$$

Our main interest is the development of robust multilevel solvers for diffusion problems

$$
\begin{equation*}
-\nabla A \nabla u=f \text { in } \Omega \subset \mathbb{R}^{2} \tag{2.4}
\end{equation*}
$$

with $A:=\left(a_{\alpha \beta}\right), a_{\alpha \beta} \in L_{\infty}(\Omega)$ and associated bilinear form

$$
\begin{equation*}
a(u, v):=\int_{\Omega}(A \nabla u, \nabla v)=\int_{\Omega} \sum_{\alpha=1}^{d} \sum_{\beta=1}^{d} a_{\alpha \beta} \partial_{\beta} u \partial_{\alpha} v . \tag{2.5}
\end{equation*}
$$

Let $\|\cdot\|_{a}$ denote the energy norm with respect to the bilinear form $a(\cdot, \cdot)$, i.e.,

$$
\|u\|_{a}:=\sqrt{a(u, u)} .
$$

The fundamental prerequisite stated in Assumption 1 for our proof is related to the discretization of (2.5) on the finest level $J$.

Assumption 1. The finest triangulation $\mathcal{T}_{J}$ consists of triangles $T$ for which the bilinear form (2.5) satisfies the local ellipticity condition

$$
\begin{equation*}
\underline{\omega}_{T, J} \sum_{\alpha=1}^{2} \xi_{\alpha}^{2} \leq \sum_{\alpha, \beta=1}^{2} a_{\alpha \beta}(x) \xi_{\alpha} \xi_{\beta} \leq \bar{\omega}_{T, J} \sum_{\alpha=1}^{2} \xi_{\alpha}^{2}, \quad \text { for all } \quad x \in T \in \mathcal{T}_{J} \tag{2.6}
\end{equation*}
$$

with weights $\underline{\omega}_{T, J}, \bar{\omega}_{T, J}$ independent of $x \in T \in \mathcal{T}_{J}$.
Since we employ a linear finite element space $\mathcal{V}_{J}$, the derivatives of $v \in \mathcal{V}_{J}$ are constant on any element $T \in \mathcal{T}_{J}$, i.e., $\nabla v=$ const, and Assumption 1 yields the equivalence

$$
\begin{equation*}
\sum_{T \in \mathcal{T}_{J}} \underline{\omega}_{T} \mu(T)\|\nabla v\|_{T}^{2} \leq a(v, v) \leq \sum_{T \in \mathcal{T}_{J}} \bar{\omega}_{T} \mu(T)\|\nabla v\|_{T}^{2} \tag{2.7}
\end{equation*}
$$

for any $v \in \mathcal{V}$ where $\mu(T)$ denotes the area of $T$ and $\|\cdot\|_{T}$ the Euclidean norm on $T \subset \mathbb{R}^{2}$.

For any coarser element $\hat{T} \in \mathcal{T}_{j}$ with $j<J$ let us now introduce average weights $\underline{\omega}_{\hat{T}}$ and $\bar{\omega}_{\hat{T}}$ based on the weights $\underline{\omega}_{T, J}$ and $\bar{\omega}_{T, J}$ from (2.6) which are assigned to the elements $T \in \mathcal{T}_{J}$ on the finest level $J$; i.e., we define

$$
\begin{equation*}
\bar{\omega}_{\hat{T}}:=\frac{1}{\mu(\hat{T})} \sum_{\substack{T \in \mathcal{T}_{J} \\ T \subset \hat{T}}} \mu(T) \bar{\omega}_{T, J} \quad \text { and } \quad \underline{\omega}_{\hat{T}}:=\frac{1}{\mu(\hat{T})} \sum_{\substack{T \in \mathcal{T}_{J} \\ T \subset \hat{T}}} \mu(T) \underline{\omega}_{T, J} \tag{2.8}
\end{equation*}
$$

Note that for elements $T \in \mathcal{T}_{J}$ on the finest level $J$ we have $\underline{\omega}_{T}=\underline{\omega}_{T, J}$ and $\bar{\omega}_{T}=\bar{\omega}_{T, J}$. Let us further introduce associated weighted norms on all levels $j$

$$
\begin{equation*}
\|v\|_{j, \bar{\omega}}^{2}:=\sum_{T \in \mathcal{T}_{j}} \bar{\omega}_{T} \int_{T}|v|^{2} \quad \text { and } \quad\|v\|_{j, \underline{\omega}}^{2}:=\sum_{T \in \mathcal{T}_{j}} \underline{\omega}_{T} \int_{T}|v|^{2} \tag{2.9}
\end{equation*}
$$

Now the fact that $\nabla v$ is constant on $T \in \mathcal{T}_{j}$ implies that (2.7) is equivalent to

$$
\begin{equation*}
\|\nabla v\|_{j, \underline{\omega}}^{2} \leq a(v, v) \leq\|\nabla v\|_{j, \bar{\omega}}^{2}, \quad \text { for } \quad j=0,1, \ldots, J \tag{2.10}
\end{equation*}
$$

In later sections we will work only with weights $\omega_{T, J}=\underline{\omega}_{T, J}=\bar{\omega}_{T, J}$ for the ease of notation. Hence, we are essentially dealing with the weighted norms $\|\cdot\|_{j, \omega}=$ $\|\cdot\|_{j, \bar{\omega}}=\|\cdot\|_{j, \underline{\omega}}$ from (2.9). Throughout the paper we also use the short hand notation

$$
\begin{equation*}
\langle u, v\rangle_{\omega}:=\sum_{T \in \mathcal{T}_{J}} \omega_{T} \int_{T} u v, \quad\|u\|_{\omega}:=\sqrt{\langle u, u\rangle_{\omega}}=\|u\|_{J, \omega} \tag{2.11}
\end{equation*}
$$

for the weighted norm on the finest level $J$.
Note that the considered $\omega_{T, J}=\underline{\omega}_{T, J}=\bar{\omega}_{T, J}$ case corresponds to the bilinear form $a(\cdot, \cdot)$ with coefficient functions $a_{\alpha \beta}=\omega \delta_{\alpha \beta}$, i.e.

$$
\begin{equation*}
a(u, v):=\int_{\Omega} \omega(\nabla u, \nabla v) \tag{2.12}
\end{equation*}
$$

and a piecewise constant approximation

$$
\omega_{T, J}:=\frac{1}{\mu(T)} \int_{T} \omega
$$

of the scalar coefficient $\omega$ with respect to the finest triangulation $\mathcal{T}_{J}$; i.e., we consider the bilinear form

$$
a_{J}(u, v)=\sum_{T \in \mathcal{T}_{J}} \omega_{T} \int_{T}(\nabla u, \nabla v)
$$

on the finest level $J$.
Remark 2.1. For the development of a preconditioner it is essential to employ identical operators and norms for both bounds, the lower and the upper bound. Therefore, it is necessary to consider the case $\omega_{T, J}=\underline{\omega}_{T, J}=\bar{\omega}_{T, J}$. Note, however, that respective upper and lower bounds can be obtained in terms of the weights $\underline{\omega}_{T}$ and $\bar{\omega}_{T}$ respectively.

For the weighted norm $\|\cdot\|_{J, \omega}=\|\cdot\|_{\omega}$ there holds a local Bernstein-type inequality under a rather weak additional condition on the weights $\omega_{T, J}$, see [17] for details.

Lemma 2.2. Consider subsets $E \subset \mathcal{T}_{J}$ of elements on the finest level $J$ which are contained in an element $T$ on a coarser level $j$, i.e. $E \subset T \in \mathcal{F}_{j}$, such that

$$
\begin{equation*}
\frac{\omega_{E}}{\omega_{T}} \leq \frac{1}{2} \frac{\mu(T)}{\mu(E)} \quad \text { and } \quad \frac{\mu(E)}{\mu(T)} \leq \gamma \tag{2.13}
\end{equation*}
$$

hold for some constant $\gamma$ independent of $E$ and $T$ where the weight $\omega_{E}$ is defined according to (2.8). Then there holds the local Bernstein-type inequality

$$
\begin{equation*}
\|\nabla v\|_{\omega, T}^{2} \leq 4 \gamma^{-2} C_{0}^{2}(\operatorname{diam}(T))^{-2}\|v\|_{\omega, T}^{2} \tag{2.14}
\end{equation*}
$$

for all $v \in \mathcal{V}_{j}$. Here $\|\cdot\|_{\omega, T}$ denotes the restriction of the norm $\|\cdot\|_{\omega}$ to $T \in \mathcal{T}_{j}$ and the constant $C_{0}$ is only dependent on the initial triangulation $\mathcal{T}_{0}$, i.e.,

$$
C_{0}:=\max _{\hat{T} \in \mathcal{T}_{0}} \operatorname{diam}(\hat{T})\left\|H_{\hat{T}}^{-1}\right\|
$$

where $H_{T}$ denotes the mapping from $T$ to the reference triangle $T_{\text {ref }}$.
Proof. Here, we only give a sketch of the proof, see [17] for details. Let $H_{T}$ denote the mapping from $T$ to the reference triangle $T_{\text {ref }}$ with vertices $(0,0),(0,1),(1,0)$. For any $v \in \mathcal{V}_{j}$ we have the representation

$$
v(x)=q(\xi)=(g, \xi)+e \quad \text { for all } x \in T, H_{T}(x)=\xi \in T_{\mathrm{ref}}
$$

and by definition

$$
\|v\|_{\omega, T}^{2}=\sum_{\substack{\hat{T} \in \mathcal{T}_{J} \\ \hat{T} \subset T}} \omega_{\hat{T}} \mu(\hat{T})\left|\left(g, \eta_{\hat{T}}\right)+e\right|^{2} \quad \text { and } \quad|\nabla v(x)|^{2} \leq C_{0}^{2}(\operatorname{diam}(T))^{-2}\|g\|^{2}
$$

for all $x \in T$, where $\eta_{\hat{T}}$ denotes the reference point in $T_{\text {ref }}$ for which $H_{T}\left(y_{\hat{T}}\right)=\eta_{\hat{T}}$ for suitable $y_{\hat{T}} \in \hat{T}$. Hence, it suffices to prove the inequality

$$
\begin{equation*}
\mu(T) \omega_{T}=\sum_{\substack{\hat{T} \in \mathcal{I}_{J} \\ \hat{T} \subset T}} \omega_{\hat{T}} \mu(\hat{T}) \leq c \sum_{\substack{\hat{T} \in \mathcal{I}_{J} \\ \hat{T} \subset T}} \omega_{\hat{T}} \mu(\hat{T}) \frac{\left|\left(g, \xi_{\hat{T}}\right)\right|^{2}}{\|g\|} \tag{2.15}
\end{equation*}
$$

We first consider the case $e=0$ and set $b:=\frac{g}{\|g\|}$. Then we can find a sub-region $S_{\text {ref }, \delta} \subset T_{\text {ref }}$ for any $\delta \in(0,1)$ such that

$$
S_{\mathrm{ref}, \delta} \subset S_{\mathrm{ref}, \delta}^{b}:=\left\{\xi \in T_{\mathrm{ref}}:|(b, \xi)| \geq \delta\right\}
$$

holds for all $b$ with $\|b\|=1$. For the set $E_{\text {ref }}:=T_{\text {ref }} \backslash S_{\text {ref }, \delta}$ one can then show that $\mu\left(E_{\text {ref }}\right) \lesssim \mu\left(T_{\text {ref }}\right)$. It is clear that the most significant contribution to the validity of $(2.15)$ is due to the points $\eta_{\hat{T}}$ which are located in $S_{\mathrm{ref}, \delta}$. With this in mind and the estimate

$$
\sum_{\substack{\hat{T} \in \mathcal{I}_{J} \\ \hat{T} \subset T}} \omega_{\hat{T}} \mu(\hat{T})\left|\left(b, \eta_{\hat{T}}\right)\right|^{2} \geq \delta^{2} \sum_{\substack{\hat{T} \in \mathcal{T}_{J} \\ \eta_{\hat{T}} \in \mathcal{S}_{\text {ref }, \delta}}} \omega_{\hat{T}} \mu(\hat{T}) \geq \delta^{2} \sum_{\substack{\hat{T} \in \mathcal{I}_{J} \\ \hat{T} \subset T}} \omega_{\hat{T}} \mu(\hat{T})\left(1-\frac{\mu(E) \omega_{E}}{\mu(T) \omega_{T}}\right)
$$

we see that we can choose $E \subset T$ such that for all $\eta_{\hat{T}} \in S_{\mathrm{ref}, \delta}$ we have $\eta_{\hat{T}} \in E$ and $\mu(E) \leq 4 \delta \mu(T)$. Hence, with a choice of $\delta=\frac{\gamma}{4}$ we obtain the asserted inequality (2.15) for $e=0$ with $c=\frac{1}{\delta^{2}}$. The case $e \neq 0$ then follows from this, see [17] for details.

Remark 2.3. Note that condition (2.13) is equivalent to the following requirement on the coefficient function $\omega$

$$
F(\gamma):=\sup _{T \in \mathcal{T}_{j}} \sup _{E \subset T, \mu(E) \leq \gamma \mu(T)} \frac{\int_{E} \omega}{\int_{T} \omega} \rightarrow 0
$$

when $\gamma \rightarrow 0$. In fact, the proof of Lemma 2.2 (see [17]) shows that it is sufficient to require $F(\gamma) \in(0,1)$. However, in this modified form condition (2.13) is equivalent to the property $A_{\infty}$ in the theory of weighted inequalities for singular integrals (see [18], p.196). It comprises the properties $A_{p}, 1 \leq p<\infty$, which give necessary and sufficient conditions in order that weighted $L_{p}$-inequalities hold for a general class of singular integrals.

With the help of Lemma 2.2 we obtain the equivalence of the weighted norms $\|\cdot\|_{j, \omega}$ and $\|\cdot\|_{\omega}$ for $v \in \mathcal{V}_{j}$ for all levels $j$.
Theorem 2.4. Under the condition (2.13) of Lemma 2.2 there holds for $v \in \mathcal{V}_{j}$

$$
\begin{equation*}
\left(1+2 \gamma^{-1} C_{0}\right)^{-1}\|v\|_{j, \omega, T} \leq\|v\|_{\omega, T} \leq\left(1+\sqrt{\frac{3}{2} C_{0}}\right)\|v\|_{j, \omega, T} \tag{2.16}
\end{equation*}
$$

with the constant $C_{0}$ only depending on the initial triangulation $\mathcal{T}_{0}$.
Proof. See [17].
This lemma is essential for our further considerations, since not all steps of the proof can be completed for both norms directly. In fact, some intermediate results can only be obtained directly either for the $\|\cdot\|_{\omega}$ norm or the $\|\cdot\|_{j, \omega}$ norms. Hence, the complete proof of the overall norm equivalencies makes significant use of this lemma.

Remark 2.5. The proof of Theorem 2.4 makes use of the fact that we are in two space dimension and employ linear finite elements. The extension of the result to higher dimensions and multi-linear elements is not obvious, however, we believe that a similar estimate holds true.

Finally, let us introduce the short hand notation $A \asymp B$ for two-sided inequalities $c A \leq B \leq C A$, and $A \succeq B, A \preceq B$ for one-sided inequalities $A \geq C B, A \leq c B$ with generic constants $c$ and $C$ which do not depend on the arguments of $A$ and $B$.

## 3. Robust Norm Equivalencies

The aim of this section is to establish two robust norm equivalencies which in turn can be used to derive robust multilevel preconditioners for our model problem (2.12). Namely, we are interested in the equivalence

$$
a(u, u) \asymp \sum_{j=1}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2}+a\left(u_{0}, u_{0}\right)
$$

and the equivalence

$$
a(u, u) \asymp \sum_{j=1}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2}+a\left(u_{0}, u_{0}\right)
$$

where the decomposition $u_{j}:=P_{j} u-P_{j-1} u$ is based on appropriate bounded and surjective projection operators $P_{j}: \mathcal{V}_{J} \rightarrow \mathcal{V}_{j}$.
3.1. Upper Bounds for the Bilinear form. We begin with the derivation of an upper bound for the bilinearform $a(u, u)$. A straightforward computation using (2.10) shows that we have the following Bernstein-type inequality relating the bilinear form (2.5) to the weighted norms from (2.9).
Lemma 3.1. For uniformly refined triangulations $\mathcal{T}_{j}$ there holds the estimate

$$
\begin{equation*}
a(v, v) \leq \frac{3}{2} C_{0}^{2} 2^{2 j}\|v\|_{j, \omega}^{2} \tag{3.1}
\end{equation*}
$$

for all $v \in \mathcal{V}_{j}$, if the bilinear form satisfies the local ellipticity condition stated in Assumption 1.

Proof. See [17], where the result is obtained for the more general weight $\bar{\omega}$.

With the help of this lemma and a strengthened Cauchy-Schwarz inequality we can derive the following sub-optimal but robust upper bound, see $[17,23]$ for details.

Theorem 3.2. Let the bilinear form $a(\cdot, \cdot)$ fulfill Assumption 1 and consider a sequence of uniformly refined triangulations $\mathcal{T}_{j}$ and the respective sequence of nested spaces $\mathcal{V}_{j}$ of linear finite elements. Then, the estimate

$$
\begin{equation*}
a(u, u) \leq 3 C_{0}^{2} J \sum_{j=1}^{J} 2^{2 j}\left\|\left(P_{j}-P_{j-1}\right) u\right\|_{j, \omega}^{2}+2 a\left(u_{0}, u_{0}\right) \tag{3.2}
\end{equation*}
$$

holds true for any choice of bounded and surjective projections $P_{j}: \mathcal{V}=\mathcal{V}_{J} \rightarrow \mathcal{V}_{j}$.
Now we will improve this sub-optimal upper estimate of $a(u, u)$. Our improvement gives a robust and optimal estimate, i.e., it does not involve the factor $J$. It is achieved in three steps: First, we introduce weighted projection operators $Q_{j}^{\omega}$ based on $\langle\cdot, \cdot\rangle_{\omega}$. Furthermore, we need to consider a second sequence of projection operators $Q_{j}^{a}$ based on auxiliary bilinear forms $a_{j}(\cdot, \cdot)$ which are defined in a leveldependent fashion. With the help of the two projections $Q_{j}^{\omega}$ and $Q_{j}^{a}$, we establish a hybrid Bernstein-type inequality involving both projections. Finally, we derive a robust and optimal upper bound of $a(\cdot, \cdot)$ using only the projections $Q_{j}^{\omega}$ via a Hardy inequality in Theorem 3.5.

To find appropriate projections for the bilinear form $a(u, v)$, we introduce the weighted projections $Q_{j}^{\omega}: \mathcal{V}_{J} \rightarrow \mathcal{V}_{j}$ by the relation

$$
\begin{equation*}
\left\langle Q_{j}^{\omega} u, v\right\rangle_{\omega}=\langle u, v\rangle_{\omega} \tag{3.3}
\end{equation*}
$$

for all $u \in \mathcal{V}_{J}$ and $v \in \mathcal{V}_{j}$. Furthermore, we need to consider auxiliary projection operators $Q_{j}^{a}: \mathcal{V}_{J} \rightarrow \mathcal{V}_{j}$ defined by

$$
\begin{equation*}
a_{j}\left(Q_{j}^{a} u, v\right)=a_{j}(u, v) \tag{3.4}
\end{equation*}
$$

for all $u \in \mathcal{V}_{J}$ and $v \in \mathcal{V}_{j}$ to obtain an optimal estimate of $a(u, u)$ by the $Q_{j}^{\omega}$. Here, the bilinear forms $a_{j}(u, v)$ are defined in a level-dependent way, i.e.,

$$
\begin{equation*}
a_{j}(u, v):=\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{T}(\nabla u, \nabla v) \tag{3.5}
\end{equation*}
$$

using the average weights (2.8) for all $u \in \mathcal{V}_{J}$ and $v \in \mathcal{V}_{j}$. Due to the definition of $\omega_{T}$ and the use of linear finite elements, it follows that

$$
\begin{aligned}
a_{k}(u, v) & =\sum_{U \in \mathcal{T}_{k}} \omega_{U} \int_{U}(\nabla u, \nabla v)=\left.\sum_{T \in \mathcal{T}_{j}} \sum_{U \in \mathcal{T}_{k}, U \subset T} \mu(U) \omega_{U}(\nabla u, \nabla v)\right|_{T} \\
& =\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{T}(\nabla u, \nabla v)=a_{j}(u, v)
\end{aligned}
$$

for $u, v \in \mathcal{V}_{j}$ with $j \leq k$. This holds true, since $\nabla u$ is constant on any $T \in \mathcal{T}_{k}$ with $k \geq j$ for $u \in \mathcal{V}_{j}$. With the definition

$$
\begin{equation*}
v_{j}=v_{j}(u):=Q_{j}^{a} u-Q_{j-1}^{a} u \tag{3.6}
\end{equation*}
$$

we obtain

$$
a\left(v_{j}, v_{k}\right)=a_{k}\left(v_{j}, Q_{k}^{a} u-Q_{k-1}^{a} u\right)=a_{k}\left(v_{j}, u\right)-a_{k}\left(v_{j}, u\right)=0
$$

for $j<k$, and since $u=Q_{J}^{a} u=\sum_{j=0}^{J} v_{j}$, where $v_{0}:=Q_{0}^{a} u$, we get the equivalence

$$
\begin{equation*}
a(u, u)=\sum_{j, k=0}^{J} a\left(v_{j}, v_{k}\right)=\sum_{j=0}^{J} a\left(v_{j}, v_{j}\right)=\sum_{j=0}^{J}\left\|v_{j}\right\|_{a}^{2} \tag{3.7}
\end{equation*}
$$

for any $u \in \mathcal{V}_{J}$. Furthermore, we introduce the sequence

$$
\begin{equation*}
u_{k}=u_{k}(u):=Q_{k}^{\omega} u-Q_{k-1}^{\omega} u \tag{3.8}
\end{equation*}
$$

based on the projections $Q_{k}^{\omega}$.
With the definitions (3.6) and (3.8) we can now obtain an upper bound for $\left\|v_{j}\right\|_{a}$ in terms of $\left\|u_{k}\right\|_{k, \omega}$ for $k \geq j$; i.e., we establish a hybrid Bernstein-type estimate of $a\left(v_{j}, v_{j}\right)$ in terms of $\left\|v_{j}\right\|_{a}\left\|u_{k}\right\|_{k, \omega}$ for $k \geq j$.

Lemma 3.3. Let $v_{j} \in \mathcal{V}_{j}$ and $u_{k} \in \mathcal{V}_{k}$ for $k \geq j$ be defined as in (3.6) and (3.8), respectively. Then there holds the estimate

$$
\begin{aligned}
a\left(v_{j}, v_{j}\right) & \leq 6 C_{2}\left\|v_{j}\right\|_{a} \sum_{k=j}^{J} 2^{2 j-k}\left\|u_{k}\right\|_{k, \omega} \\
\left\|v_{j}\right\|_{a} & \leq 6 C_{2} \sum_{k=j}^{J} 2^{2 j-k}\left\|u_{k}\right\|_{k, \omega},
\end{aligned}
$$

with a constant

$$
\begin{equation*}
C_{2}:=\max _{\hat{T} \in \mathcal{T}_{0}} \frac{\operatorname{diam}(\hat{T})}{\sqrt{\mu(\hat{T})}} \tag{3.9}
\end{equation*}
$$

which depends on the initial triangulation $\mathcal{T}_{0}$ only.
Proof. Since $v_{j}=Q_{j}^{a} u-Q_{j-1}^{a} u$, we have $a\left(v_{j}, w\right)=0$ for all $w \in \mathcal{V}_{j-1}$. Hence, with the choice $w=Q_{j-1}^{\omega} u$ we obtain

$$
\begin{aligned}
a\left(v_{j}, v_{j}\right) & =a_{j}\left(v_{j}, v_{j}\right)=a_{j}\left(v_{j}, u\right)=a_{j}\left(v_{j}, u-Q_{j-1}^{\omega} u\right) \\
& =\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{T}\left(\nabla v_{j}, \nabla\left(u-Q_{j-1}^{\omega} u\right)\right) .
\end{aligned}
$$

Using the identity $u-Q_{j-1}^{\omega} u=\sum_{k=j}^{J} Q_{k}^{\omega} u-Q_{k-1}^{\omega} u=\sum_{k=j}^{J} u_{k}$ we establish the equivalence

$$
a\left(v_{j}, v_{j}\right)=\sum_{T \in \mathcal{T}_{j}} \omega_{T} \sum_{k=j}^{J} \int_{T}\left(\nabla v_{j}, \nabla u_{k}\right)
$$

and integration by parts for each $T \in \mathcal{T}_{j}$ yields

$$
a\left(v_{j}, v_{j}\right)=\sum_{T \in \mathcal{T}_{j}} \sum_{k=j}^{J} \omega_{T} \int_{\partial T} u_{k}\left(\nabla v_{j}, n_{\partial T}\right) .
$$

Now consider a fixed $T \in \mathcal{T}_{j}$ and $U \in \mathcal{T}_{k}$ with $U \subset T$. Then we obtain

$$
\begin{aligned}
\left|\omega_{T} \int_{\partial T} u_{k}\left(\nabla v_{j}, n_{\partial T}\right)\right| & =\left|\mu(T)^{-1} \sum_{\substack{U \in \mathcal{I}_{k} \\
U \subset T}} \mu(U) \omega_{U} \int_{\partial U} u_{k}\left(\nabla v_{j}, n_{\partial U}\right)\right| \\
& \leq \mu(T)^{-1} \sum_{\substack{U \in \mathcal{I}_{k} \\
U \subset T}} \mu(U) \sqrt{\omega_{U}} \int_{\partial U}\left|u_{k}\right|\left|\nabla v_{j}\right|_{\infty, T} \sqrt{\omega_{U}} \\
& \leq 6 \sum_{\substack{U \in \mathcal{I}_{k} \\
U \subset T}} \sqrt{\omega_{U} \int_{U}\left|u_{k}\right|^{2}} \frac{\operatorname{diam}(U)}{\mu(T)} \sqrt{\omega_{U} \int_{U}\left|\nabla v_{j}\right|^{2}}
\end{aligned}
$$

Here, in the last line, the inequality

$$
\int_{\partial U}\left|u_{k}\right| \leq \operatorname{diam}(U)\left[w_{1}+w_{2}+w_{3}\right] \leq 6 \operatorname{diam}(U) \mu(U)^{-\frac{1}{2}}\left(\int_{U}\left|u_{k}\right|^{2}\right)^{\frac{1}{2}}
$$

was used which results from the well-known formula

$$
\int_{U}|v|^{2}=\frac{\mu(U)}{12}\left[w_{1}^{2}+w_{2}^{2}+w_{3}^{2}+\left(w_{1}+w_{2}+w_{3}\right)^{2}\right]
$$

for linear functions $v$ on $U$ with vertices $w_{1}, w_{2}$, and $w_{3}$. Since we consider uniformly refined triangulations in two dimensions, we have $\mu(T)=2^{2(k-j)} \mu(U)$ and therefore $\operatorname{diam}(U) \leq C_{2} 2^{2(j-k)} 2^{k} \mu(T)$ with $C_{2}$ given in (3.9) depending on the initial triangulation $\mathcal{T}_{0}$ only. Hence, we end up with the assertion

$$
\begin{aligned}
a\left(v_{j}, v_{j}\right) & \leq 6 C_{2} \sum_{k=j}^{J} 2^{-k} \sum_{T \in \mathcal{T}_{j}}\left\|u_{k}\right\|_{k, \omega, T} 2^{2 j}\left\|\nabla v_{j}\right\|_{j, \omega, T} \\
& \leq 6 C_{2} \sum_{k=j}^{J} 2^{2 j-k}\left\|u_{k}\right\|_{k, \omega}\left\|v_{j}\right\|_{a}
\end{aligned}
$$

Note that this lemma is a strengthened version of the following Bernstein-type inequality of broken order for certain weighted trace norms

$$
\|u\|_{\frac{1}{2}, j, \omega}:=\left(\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{\partial T}|u|^{2}\right)^{\frac{1}{2}}
$$

which employs only a single projection.
Lemma 3.4. For elements $v \in \mathcal{V}_{j}$ there holds

$$
\|v\|_{a} \leq \sqrt{3 C_{2} C_{1}} 2^{\frac{j}{2}}\|v\|_{\frac{1}{2}, j, \omega}
$$

where the constants $C_{1}:=\max _{T \in \mathcal{T}_{0}} \operatorname{diam}(T)$, and $C_{2}$ from (3.9) depend on the initial triangulation $\mathcal{T}_{0}$ only.

Proof. Keeping in mind that $\nabla v$ is constant on $T \in \mathcal{T}_{j}$, integration by parts yields

$$
\begin{aligned}
a(v, v) & =\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{T}(\nabla v, \nabla v)=\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{\partial T} v\left(\nabla v, n_{\partial T}\right) \\
& \leq\left(\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{\partial T}|v|^{2}\right)^{\frac{1}{2}}\left(\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{\partial T}\|\nabla v\|^{2}\right)^{\frac{1}{2}}
\end{aligned}
$$

Here, the first sum represents the semi-norm $\|v\|_{\frac{1}{2}, \omega, j}$, whereas each term in the second sum can be bounded using the local Bernstein-type inequality

$$
\int_{\partial T}\|\nabla v\|^{2} \leq 3 \operatorname{diam}(T)\|\nabla v\|^{2} \leq 3 C_{2} \mu(T)^{\frac{1}{2}}\|\nabla v\|^{2} \leq 3 C_{2} C_{1} 2^{j} \int_{T}\|\nabla v\|^{2}
$$

Hence, after multiplication with $\omega_{T}$ and summation with respect to $T \in \mathcal{T}_{j}$, we obtain the overall estimate

$$
a(v, v)=\sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{T}(\nabla v, \nabla v) \leq \sqrt{3 C_{2} C_{1}} 2^{\frac{j}{2}}\|v\|_{\frac{1}{2}, j, \omega} \sqrt{a(v, v)}
$$

and the assertion follows after division by $\sqrt{a(v, v)}$.
Finally, we are in the position to prove the main result of this section, a robust and optimal upper bound for the bilinear form $a(u, u)$.

Theorem 3.5. Let the bilinear form $a(\cdot, \cdot)$ fulfill Assumption 1 and consider a sequence of uniformly refined triangulations $\mathcal{T}_{j}$ and the respective sequence of nested spaces $\mathcal{V}_{j}$ of linear finite elements. Then, there holds the upper bound

$$
\begin{equation*}
a(u, u) \leq \frac{32}{3} C_{2} \sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2} \tag{3.10}
\end{equation*}
$$

for $u_{k}$ from (3.8) and the norm $\|\cdot\|_{j, \omega}$ from (2.9). The constant $C_{2}$ given in (3.9) depends only on the initial triangulation $\mathcal{T}_{0}$.
Proof. With (3.7) and Lemma 3.3 we obtain

$$
\begin{align*}
a(u, u) & =\sum_{j=0}^{J}\left\|v_{j}\right\|_{a}^{2} \leq 6 C_{2} \sum_{k=0}^{J} 2^{-k}\left\|u_{k}\right\|_{k, \omega} \sum_{j=0}^{k} 2^{2 j}\left\|v_{j}\right\|_{a} \\
& \leq\left(\sum_{k=0}^{J} 2^{2 k}\left\|u_{k}\right\|_{k, \omega}^{2}\right)^{\frac{1}{2}}\left(\sum_{k=0}^{J} 16^{-k}\left(\sum_{j=0}^{k} 2^{2 j}\left\|v_{j}\right\|_{a}\right)^{2}\right)^{\frac{1}{2}} \tag{3.11}
\end{align*}
$$

after interchanging the sums. Using the Hardy inequality

$$
\begin{equation*}
\left(\sum_{k=0}^{J} b^{k} s_{k}^{2}\right)^{\frac{1}{2}} \leq \frac{1}{1-\sqrt{b}}\left(\sum_{j=0}^{J} b^{j} a_{j}^{2}\right)^{\frac{1}{2}} \tag{3.12}
\end{equation*}
$$

where

$$
s_{k}:=\sum_{j=0}^{k} a_{j}, \quad s_{-1}:=0, \quad \text { and } \quad c_{k}:=\sum_{l=k}^{J} b^{l}, \quad c_{J+1}:=0
$$

we obtain

$$
\left(\sum_{k=0}^{J} 16^{-k}\left(\sum_{j=0}^{k} 2^{2 j}\left\|v_{j}\right\|_{a}\right)^{2}\right)^{\frac{1}{2}} \leq \frac{4}{3}\left(\sum_{j=0}^{J}\left\|v_{j}\right\|_{j, \omega}^{2}\right)^{\frac{1}{2}}
$$

with the choice $a_{j}:=2^{2 j}\left\|v_{j}\right\|_{j, \omega}$ and $b=1 / 16$. Plugging this estimate into (3.11), we establish the asserted optimal and robust upper bound (3.10).

We obtain the respective upper bound for the weighted norm $\|\cdot\|_{\omega}$ on the finest level using Theorem 2.4.

Theorem 3.6. Let the bilinear form $a(\cdot, \cdot)$ fulfill Assumption 1 and consider a sequence of uniformly refined triangulations $\mathcal{T}_{j}$ and the respective sequence of nested spaces $\mathcal{V}_{j}$ of linear finite elements. Then, there holds the upper bound

$$
\begin{equation*}
a(u, u) \leq\left(1+2 \gamma^{-1} C_{0}\right)^{2} \frac{32}{3} C_{2} \sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \tag{3.13}
\end{equation*}
$$

for $u_{j}$ given in (3.8) and the norm $\|\cdot\|_{\omega}$ from (2.11).
3.2. Lower Bounds for the Bilinear form. The next step in our search for robust norm equivalencies is the derivation of optimal and robust lower bounds for the bilinear form $a(\cdot, \cdot)$ in terms of the norms $\|\cdot\|_{j, \omega}$ and $\|\cdot\|_{\omega}$; i.e., we are looking for a Jackson-type inequality for $\left\|u_{j}\right\|_{j, \omega}$ and for $\left\|u_{j}\right\|_{\omega}$
$\sum_{j=1}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2}+a\left(u_{0}, u_{0}\right) \preceq a(u, u) \quad$ and $\quad \sum_{j=1}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2}+a\left(u_{0}, u_{0}\right) \preceq a(u, u)$.
Unlike in the Bernstein case, the constants will involve some information about the coefficient $\omega$, yet in a very weak norm. The main reason for this slight dependence is due to the problem that we cannot obtain an optimal lower bound directly for the $\|\cdot\|_{j, \omega}$ norms. Here, we rather need to work with the $\|\cdot\|_{\omega}$ norm to obtain an optimal estimate. Then, we can exploit the result of Theorem 2.4 to get a similar bound for $\|\cdot\|_{j, \omega}$. Yet, it is this detour in our proof, which forces us to introduce some information about the weight $\omega$ in the estimate.

In a first step, we bound the $u_{j}$-decomposition in terms of the $v_{j}$-decomposition with respect to the $\|\cdot\|_{\omega}$ norm in the following lemma. The respective estimate for the $\|\cdot\|_{j, \omega}$ norms then follows with Theorem 2.4.
Lemma 3.7. The decompositions defined in (3.6) and (3.8) allow for the estimates

$$
\sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \leq 4 \sum_{j=0}^{J} 2^{2 j}\left\|v_{j}\right\|_{\omega}^{2}
$$

and

$$
\sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2} \leq 4\left(1+2 \gamma^{-1} C_{0}\right)^{2}\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} \sum_{j=0}^{J} 2^{2 j}\left\|v_{j}\right\|_{j, \omega}^{2}
$$

Proof. Observe that due to (3.3) we have

$$
\begin{aligned}
&\left\|u_{j}\right\|_{\omega}^{2}=\left\|Q_{j}^{\omega} u-Q_{j-1}^{\omega} u\right\|_{\omega}^{2} \\
&=\left\langle Q_{j}^{\omega} u-Q_{j-1}^{\omega} u, u\right\rangle_{\omega}=\left\langle Q_{j}^{\omega} u-Q_{j-1}^{\omega} u, Q_{j}^{\omega} u\right\rangle_{\omega} \\
&=\left\langle Q_{j-1}^{\omega} u, u-Q_{j-1}^{a} u\right\rangle_{\omega}
\end{aligned}
$$

for all $j$ so that we obtain

$$
\sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \leq \sum_{j=0}^{J} 2^{2 j}\left\|u-Q_{j-1}^{a} u\right\|_{\omega}^{2} \leq \sum_{j=0}^{J} 2^{2 j} \sum_{k=j}^{J}\left\|v_{k}\right\|_{\omega}^{2}
$$

where $v_{j}$ is defined in (3.6) and $u_{j}$ is defined in (3.8). Then using Hardy's inequality again, we obtain

$$
\sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \leq 4 \sum_{j=0}^{J} 2^{2 j}\left\|v_{j}\right\|_{\omega}^{2}
$$

Now we pass back from the $\|\cdot\|_{\omega}$ norm to the $\|\cdot\|_{j, \omega}$ norms with the help of Theorem 2.4 and end up with the estimates

$$
\begin{equation*}
\sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \leq 4\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} \sum_{j=0}^{J} 2^{2 j}\left\|v_{j}\right\|_{j, \omega}^{2} \tag{3.14}
\end{equation*}
$$

and

$$
\sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2} \leq 4\left(1+2 \gamma^{-1} C_{0}\right)^{2}\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} \sum_{j=0}^{J} 2^{2 j}\left\|v_{j}\right\|_{j, \omega}^{2}
$$

Remark 3.8. Note that it is necessary for the proof to switch to the $\|\cdot\|_{\omega}$ norm, since the terms in the Hardy inequality must employ the same norm on all levels $j$.

Now, we need to deal with the projection operators $Q_{j}^{a}$ and the respective sequence $v_{j}$ only. Here, we can prove a local estimate for the $\|\cdot\|_{j, \omega}$ norms using a modified version of the duality technique due to Aubin and Nitsche, see also [2], and a certain regularity result for the Neumann problem. However, we are able to obtain this local estimate on a particular level $j$ only under the additional assumption

$$
\begin{equation*}
1>C \max _{\substack{S^{\prime}, \in \in \mathcal{T}_{j} \\ S^{\prime}, S \in U}} \sqrt{\frac{\omega_{S}}{\omega_{S^{\prime}}}} \operatorname{diam}(U) \tag{3.15}
\end{equation*}
$$

where $C$ denotes an absolute constant, and $U=U_{T}$ refers to the union of the supports of all nodal basis functions $\psi_{l, j-1}$ which intersect the considered element $T \in \mathcal{T}_{j-1}$ (compare (3.16)). Note that this condition essentially intertwines the necessary refinement level $j$ with the "smoothness" of the coefficient function $\omega$.

Theorem 3.9. Consider the neighborhood

$$
\begin{equation*}
U:=U_{T}=\bigcup_{\substack{\hat{T} \in \mathcal{T}_{j-1} \\ \hat{T} \subset \operatorname{supp}\left(\psi_{l, j-1}\right) \\ \operatorname{supp}\left(\psi_{l, j-1}\right) \cap T \neq \emptyset}} \hat{T} \tag{3.16}
\end{equation*}
$$

of $T \in \mathcal{T}_{j-1}$ and let condition (3.15) be fulfilled. Then there holds the estimate

$$
\begin{equation*}
\left\|v_{j}\right\|_{j, \omega, U} \leq C\left(\max _{\substack{S^{\prime}, S \in \mathcal{I}_{j} \\ S^{\prime}, S \subset U}} \sqrt{\frac{\omega_{S}}{\omega_{S^{\prime}}}}\right) \operatorname{diam}(U)\left\|\nabla v_{j}\right\|_{j, \omega, U} \tag{3.17}
\end{equation*}
$$

Essential for the proof of (3.17) is the regularity result stated in the following lemma.
Lemma 3.10. Consider the domain $\tilde{U}$ with a smooth boundary $\partial \tilde{U}$ such that $\tilde{U} \supset$ $U$, all vertices $\xi \in \partial U$ are also $\xi \in \partial \tilde{U}$, and $\operatorname{diam}(\tilde{U})<2 \operatorname{diam}(U)$ with $U$ defined in (3.16). Let $\tilde{v}_{j}$ denote the continuous and piecewise linear extension of $v_{j}$ from $U$ to $\tilde{U}$ which fulfills

$$
\begin{equation*}
\left\|\tilde{v}_{j}\right\|_{0,2, \tilde{U}} \leq 2\left\|v_{j}\right\|_{0,2, U} \tag{3.18}
\end{equation*}
$$

Then, the solution $\varphi_{U}$ to the inhomogeneous Neumann problem

$$
-\Delta \varphi_{U}=\tilde{v}_{j} \text { in } \tilde{U}, \quad \frac{\partial \varphi_{U}}{\partial \nu}=g \text { on } \partial \tilde{U}
$$

is in $H^{2}(U)$ and allows for the estimate

$$
\begin{equation*}
\left|\varphi_{U}\right|_{2,2, U} \leq C\left\|v_{j}\right\|_{0,2, U} \tag{3.19}
\end{equation*}
$$

Here, the boundary data for the Neumann problem are $g:=\alpha h$ with a piecewise linear $h \in L_{2}(\partial \tilde{U})$ such that $h \geq 0$ on the boundary $\partial \tilde{U}$ and $h(\xi)=0$ for all vertices $\xi \in \partial U \cap \partial \tilde{U}$. Furthermore, $\alpha \in \mathbb{R}$ is such that the compatibility condition

$$
\begin{equation*}
\int_{\tilde{U}} \tilde{v}_{j}=-\int_{\partial \tilde{U}} g \tag{3.20}
\end{equation*}
$$

is fulfilled.
Proof. Consider the scaled domain $\hat{\Omega}:=R \tilde{U}$ and the spaces

$$
\begin{aligned}
V(\hat{\Omega}) & :=\left\{\varphi \in W_{2}^{2}(\hat{\Omega}): \int_{\hat{\Omega}} \varphi d x=0\right\} \quad \text { and } \\
W(\hat{\Omega}) & :=\left\{\langle f, g\rangle \in L_{2}(\hat{\Omega}) \times W_{2}^{1 / 2}(\partial \hat{\Omega}): \int_{\hat{\Omega}} f d x+\int_{\partial \hat{\Omega}} g d s=0\right\}
\end{aligned}
$$

and the mapping $T: \varphi \in V \mapsto T \varphi \in W$ defined by

$$
T \varphi:= \begin{cases}f=\Delta \varphi & \in L_{2}(\hat{\Omega}) \\ g=\partial \varphi / \partial \nu & \in W_{2}^{1 / 2}(\partial \hat{\Omega})\end{cases}
$$

This mapping is linear and continuous; i.e., there exists a constant $M_{\hat{\Omega}}$ such that

$$
\|T \varphi\|_{W} \equiv\|f\|_{0,2, \hat{\Omega}}+\|g\|_{1 / 2,2, \partial \hat{\Omega}} \leq M_{\hat{\Omega}}\|\varphi\|_{2,2, \hat{\Omega}}
$$

holds. Furthermore, $T$ is also bijective from $V$ onto $W$, see [21], p. 336-339. Hence by the open mapping theorem its inverse $T^{-1}:\langle f, g\rangle \mapsto \varphi$ is also continuous and satisfies

$$
\|\varphi\|_{V} \equiv\|\varphi\|_{2,2, \hat{\Omega}} \leq L_{\hat{\Omega}}\left(\|\Delta \varphi\|_{0,2, \hat{\Omega}}+\|\nabla \varphi / \partial \nu\|_{1 / 2,2, \partial \hat{\Omega}}\right), \quad \varphi \in W(\hat{\Omega})
$$

with a constant $L_{\hat{\Omega}}$. However, it is still necessary to determine the dependence of $L_{\hat{\Omega}}$ on the size of $\hat{\Omega}=R \tilde{U}$, i.e., on the scaling $R$, since by definition (3.16) $U=U_{T}$ depends on the level $j$. To this end, let us consider the scaling $R$ such that $R^{-1}:=\operatorname{diam}(\tilde{U}) \leq 2 \operatorname{diam}(U)$ so that

$$
\begin{equation*}
\|\psi\|_{2,2, \hat{\Omega}} \leq C^{*}\left(\|\Delta \psi\|_{0,2, \hat{\Omega}}+\|\partial \psi / \partial \nu\|_{1 / 2,2, \partial \hat{\Omega}}\right), \quad \psi \in W(\hat{\Omega}) \tag{3.21}
\end{equation*}
$$

with $C^{*}$ depending only on the shape of $U$ but not on the size. Hence, $C^{*}$ depends only on the initial triangulation $\mathcal{T}_{0}$. The connection between $\varphi$ and $\psi$ is given by $\psi(t)=\varphi(t / R):=\varphi(x)$. Therefore, there hold the equivalencies

$$
\|\psi\|_{0,2, \hat{\Omega}}=R\|\varphi\|_{0,2, \tilde{U}}, \quad|\psi|_{1,2, \hat{\Omega}}=|\varphi|_{1,2, \tilde{U}}, \quad \text { and } \quad|\psi|_{2,2, \hat{\Omega}}=R^{-1}|\varphi|_{2,2, \tilde{U}}
$$

Using the explicit form of the trace norm $\|\cdot\|_{1 / 2,2, \partial \hat{\Omega}}$ given by (see e.g. [21], p. 94)

$$
\begin{align*}
\|\psi\|_{1 / 2,2, \partial \hat{\Omega}}^{2} & :=\sum_{i}\left\|\psi_{i}\right\|_{0,2, \partial \hat{\Omega}}^{2}+\sum_{i}\left|\psi_{i}\right|_{1 / 2,2, \partial \hat{\Omega}}^{2} \\
\left|\psi_{i}\right|_{1 / 2,2, \partial \hat{\Omega}}^{2} & :=\int_{\partial \hat{\Omega}} \int_{\partial \hat{\Omega}} \frac{\left|\psi_{i}(t)-\psi_{i}(s)\right|^{2}}{|t-s|^{2}} d \sigma d \sigma \tag{3.22}
\end{align*}
$$

where $\psi=\sum_{i} \psi_{i}$ is a partition of $\psi$ with respect to the representation by charts of the curve $\partial \hat{\Omega}$ with curve element $d \sigma$, we conclude

$$
\|\partial \psi / \partial \nu\|_{1 / 2,2, \partial \hat{\Omega}}=R^{-1 / 2}\|\partial \varphi / \partial \nu\|_{0,2, \partial \tilde{U}}+R^{-1}|\partial \varphi / \partial \nu|_{1 / 2,2, \partial \tilde{U}}
$$

Consequently, by (3.21) we obtain the estimate

$$
\begin{align*}
\left|\varphi_{U}\right|_{1,2, \tilde{U}}+\operatorname{diam}(U)\left|\varphi_{U}\right|_{2,2, \tilde{U}} \leq C^{*} \operatorname{diam}(U) & \left(\left\|\tilde{v}_{j}\right\|_{0,2, \tilde{U}}+\right. \\
& \left.|g|_{1 / 2,2, \partial \tilde{U}}+\frac{\|g\|_{0,2, \partial \tilde{U}}}{\sqrt{\operatorname{diam}(U)}}\right) \tag{3.23}
\end{align*}
$$

for the data of the Neumann problem above, i.e., $\varphi=\varphi_{U} \in W_{2}^{2}(\tilde{U}),-\Delta \varphi_{U}=\tilde{v}_{j}$ in $\tilde{U}$ and $\partial \varphi / \partial \nu=g$ on $\partial \tilde{U}$. Now we use the fact that $g$ is piecewise linear on $\partial \tilde{U}$ by construction. Therefore,

$$
\begin{align*}
|g|_{1 / 2,2, \partial \tilde{U}} & \leq C(\operatorname{diam}(U))^{-1}\|g\|_{0,1, \partial \tilde{U}}  \tag{3.24}\\
\|g\|_{0,2, \partial \tilde{U}} & \leq C(\operatorname{diam}(U))^{-1 / 2}\|g\|_{0,1, \partial \tilde{U}} \tag{3.25}
\end{align*}
$$

holds with a constant $C$ independent of $\tilde{U}$ and $g$. For completeness, we give the proof of the inequalities (3.24) and (3.25) here. To this end, let $\partial \tilde{U}_{i}$ be one of the two pieces of the segment of $\partial \tilde{U}$ between two vertices of $\mathcal{T}_{j-1}$. Here $g$ is a linear function. Then by (3.22) with $\psi_{i}=g$

$$
|g|_{1 / 2,2, \partial \tilde{U}_{i}}=\left(\int_{\partial \tilde{U}_{i}} \int_{\partial \tilde{U}_{i}} \frac{|g(t)-g(s)|^{2}}{|t-s|^{2}} d \sigma d \sigma\right)^{1 / 2} \leq C \operatorname{diam}\left(\partial \tilde{U}_{i}\right)|\nabla g|_{\infty, \partial \tilde{U}_{i}}
$$

follows. Furthermore, we can replace $g$ by $\tilde{g}:=g-\bar{g}$ where $\bar{g}:=\min _{\partial \tilde{U}_{i}}|g|$. Then $\tilde{g}$ has a zero on $\partial \tilde{U}_{i}$ and we have the equivalence

$$
\int_{\partial \tilde{U}_{i}}|\tilde{g}| d \sigma \asymp\left(\operatorname{diam}\left(\partial \tilde{U}_{i}\right)\right)^{2}|\nabla g|_{\infty, \partial \tilde{U}_{i}}
$$

with constants independent of $\tilde{U}_{i}$ and $g$. Since

$$
\int_{\partial \tilde{U}_{i}}|\tilde{g}| d \sigma \leq \int_{\partial \tilde{U}_{i}}|g| d \sigma+C \operatorname{diam}\left(\partial \tilde{U}_{i}\right) \bar{g} \leq C \int_{\partial U_{i}}|g| d \sigma=C\|g\|_{0,1, \partial \tilde{U}_{i}}
$$

the comparison with the previous inequality yields (3.24). The second inequality (3.25) follows from

$$
\begin{aligned}
\int_{\partial \tilde{U}_{i}}|g|^{2} d \sigma & \leq C\left(\int_{\partial \tilde{U}_{i}}|\tilde{g}|^{2} d \sigma+\int_{\partial \tilde{U}_{i}}|\bar{g}|^{2} d \sigma\right) \\
& \leq C\left(\operatorname{diam}\left(\partial \tilde{U}_{i}\right)\right)^{3}|\nabla g|_{\infty, \partial \tilde{U}_{i}}^{2}+C\left(\operatorname{diam}\left(\partial \tilde{U}_{i}\right)\right)^{-1}\left(\int_{\partial \tilde{U}_{i}}|g| d \sigma\right)^{2} \\
& \leq C\left(\operatorname{diam}\left(\partial \tilde{U}_{i}\right)\right)^{-1}\left(\int_{\partial \tilde{U}_{i}}|g| d \sigma\right)^{2} .
\end{aligned}
$$

This completes the proof of (3.24) and (3.25), since we can choose $\tilde{U}$ such that $\operatorname{diam}(\tilde{U}) \asymp \operatorname{diam}(\partial \tilde{U})$.

Inserting (3.24) and (3.25) into inequality (3.23) and taking (3.18) and (3.20) into account, the desired inequality (3.19) directly follows:

$$
\begin{aligned}
\left|\varphi_{U}\right|_{2,2, U} & \leq C\left(\left\|\tilde{v}_{j}\right\|_{0,2, \tilde{U}}+(\operatorname{diam}(U))^{-1}\|g\|_{0,1, \partial \tilde{U}}\right) \\
& =C\left(\left\|\tilde{v}_{j}\right\|_{0,2, \tilde{U}}+(\operatorname{diam}(U))^{-1}\left\|\tilde{v}_{j}\right\|_{0,1, \tilde{U}}\right) \leq C\left\|v_{j}\right\|_{0,2, U}
\end{aligned}
$$

With this lemma, we are in the position to prove Theorem 3.9.

Proof (Theorem 3.9). At first we note that we can assume $v_{j}(x) \geq 0$ in $U$ since otherwise $v_{j}$ changes sign in $U$ and we can obtain the desired inequality (3.17) directly. Therefore we need to consider in the following only $U$ which do not intersect boundary elements $T \in \mathcal{T}_{j-1}$.

With the help of (2.9) and Lemma 3.10 we obtain, after integration by parts on each $S \subset U$,
$\left\|v_{j}\right\|_{j, \omega, U}^{2}=-\int_{U} v_{j} \Delta \varphi_{U}=\sum_{S \subset U} \omega_{S} \int_{S}\left(\nabla v_{j}, \nabla \varphi_{U}\right)-\sum_{S \subset U} \omega_{S} \int_{\partial S} v_{j}\left(\nabla \varphi_{U}, n_{\partial S}\right)$.
Concerning the first sum, observe that by definition of $v_{j}$, see (3.6),

$$
0=a\left(v_{j}, w\right)=\sum_{S \in \mathcal{T}_{j}} \omega_{S} \int_{S}\left(\nabla v_{j}, \nabla w\right)
$$

holds for all $w \in \mathcal{V}_{j-1}$. Now, we choose $w$ to be the function in $\mathcal{V}_{j-1}$ which interpolates $\varphi_{U}$ at the nodes in $T$ and has support in $U=U_{T}$. Then, we can estimate the first sum via the Bramble-Hilbert Lemma on $U$

$$
\begin{align*}
I_{1}(U):=\left|\sum_{S \subset U} \omega_{S} \int_{S}\left(\nabla \varphi_{U}, \nabla v_{j}\right)\right| & =\left|\sum_{S \subset U} \omega_{S} \int_{S}\left(\nabla\left(\varphi_{U}-w\right), \nabla v_{j}\right)\right| \\
& \leq\left(\sum_{S \subset U} \omega_{S} \int_{S}\left|\nabla\left(\varphi_{U}-w\right)\right|^{2}\right)^{1 / 2}\left\|\nabla v_{j}\right\|_{j, \omega, U} \\
& \leq C \operatorname{diam}(U)\left|\varphi_{U}\right|_{2,2, U}\left(\max _{S \subset U} \sqrt{\omega_{S}}\right)\left\|\nabla v_{j}\right\|_{j, \omega, U} \tag{3.27}
\end{align*}
$$

where the constant $C$ depends only on the shape of $U$, i.e., by assumption only on the initial triangulation $\mathcal{T}_{0}$.

The second sum in (3.26) can be estimated by

$$
\left|\sum_{S \subset U} \omega_{S} \int_{\partial S} v_{j}\left(\nabla \varphi_{U}, n_{\partial S}\right)\right| \leq \sum_{S \subset U} \omega_{S}\left\|v_{j}\right\|_{\infty, \partial S} \int_{\partial S}\left|\nabla \varphi_{U}\right|:=I_{2}(U)
$$

Note that due to the choice of $g$ in Lemma 3.10, the normal derivative $\partial \varphi_{U} / \partial \nu_{\xi}=g$ vanishes at each vertex $\xi \in \partial U$. And since any normal vector $n_{K}$ of an arbitrary edge $K \subset U$ can be represented as a linear combination of three such normal vectors at $\xi_{i}, i=1,2,3$, we can bound the normal derivative at $x \in K$ by

$$
\begin{aligned}
\left|\left(\nabla \varphi_{U}, n_{K}\right)(x)\right| & =\left|\sum_{i=1}^{3} \beta_{i} \frac{\partial \varphi_{U}}{\partial \nu_{\xi_{i}}}(x)\right| \\
& =\left|\sum_{i=1}^{3} \beta_{i}\left(\frac{\partial \varphi_{U}}{\partial \nu_{\xi_{i}}}(x)-\frac{\partial \varphi_{U}}{\partial \nu_{\xi_{i}}}\left(\xi_{i}\right)\right)\right| \\
& \leq \sum_{i=1}^{3}\left|\beta_{i}\right| \int_{0}^{1}\left|\left(\nabla \frac{\partial \varphi_{U}}{\partial \nu}\right)\left(\xi_{i}+t\left(x-\xi_{i}\right), x-\xi_{i}\right) d t\right|
\end{aligned}
$$

with $\sum_{i=1}^{3}\left|\beta_{i}\right|^{2}=1$. Hence after integration over $K$ we obtain

$$
\int_{K}\left|\left(\nabla \varphi_{U}, n_{K}\right)(x)\right| d x \leq C \operatorname{diam}(U) \sum_{|\alpha|=2} \int_{U}\left|D^{\alpha} \varphi_{U}\right| d x \leq C(\operatorname{diam}(U))^{2}\left|\varphi_{U}\right|_{2,2, U}
$$

Altogether, we can now establish the estimate

$$
\begin{align*}
I_{2}(U) & \leq C \sum_{S \subset U} \omega_{S}\left\|v_{j}\right\|_{\infty, \partial S}(\operatorname{diam}(U))^{2}\left|\varphi_{U}\right|_{2,2, U} \\
& \leq C\left\|v_{j}\right\|_{j, \omega, U}\left(\max _{S \subset U} \sqrt{\omega_{S}}\right) \operatorname{diam}(U)\left|\varphi_{U}\right|_{2,2, U}, \tag{3.28}
\end{align*}
$$

due to $\sum_{S \subset U} \sqrt{\omega_{S}} \operatorname{diam}(U)\left\|v_{j}\right\|_{\infty, \partial S} \leq C\left\|v_{j}\right\|_{j, \omega, U}$. Now the assertion (3.17) follows easily with the aid of (3.19): Insert (3.19) into (3.27) and (3.28) and obtain

$$
\begin{aligned}
\left\|v_{j}\right\|_{j, \omega, U}^{2} & \leq I_{1}(U)+I_{2}(U) \\
& \leq C \operatorname{diam}(U)\left(\left\|v_{j}\right\|_{j, \omega, U}+\left\|\nabla v_{j}\right\|_{j, \omega, U}\right)\left(\max _{S \subset U} \sqrt{\omega_{S}}\right)\left\|v_{j}\right\|_{0,2, U}
\end{aligned}
$$

With the estimate $\left(\max _{S \subset U} \sqrt{\omega_{S}}\right)\left\|v_{j}\right\|_{0,2, U} \leq\left(\max _{S^{\prime}, S \subset U} \sqrt{\omega_{S} / \omega_{S^{\prime}}}\right)\left\|v_{j}\right\|_{j, \omega, U}$, this yields the inequality

$$
\left\|v_{j}\right\|_{j, \omega, U} \leq C \operatorname{diam}(U)\left(\max _{S^{\prime}, S \subset U} \sqrt{\frac{\omega_{S}}{\omega_{S^{\prime}}}}\right)\left(\left\|v_{j}\right\|_{j, \omega, U}+\left\|\nabla v_{j}\right\|_{j, \omega, U}\right)
$$

after division by $\left\|v_{j}\right\|_{j, \omega, U}$. Hence, we end up with $\left(1-C \operatorname{diam}(U)\left(\max _{S^{\prime}, S \subset U} \sqrt{\frac{\omega_{S}}{\omega_{S^{\prime}}}}\right)\right)\left\|v_{j}\right\|_{j, \omega, U} \leq C \operatorname{diam}(U)\left(\max _{S^{\prime}, S \subset U} \sqrt{\frac{\omega_{S}}{\omega_{S^{\prime}}}}\right)\left\|\nabla v_{j}\right\|_{j, \omega, U}$.

Thus, the assertion follows with assumption (3.15)

$$
\left\|v_{j}\right\|_{j, \omega, U} \leq \frac{C \operatorname{diam}(U)\left(\max _{S^{\prime}, S \subset U} \sqrt{\frac{\omega_{S}}{\omega_{S^{\prime}}}}\right)}{1-C \operatorname{diam}(U)\left(\max _{S^{\prime}, S \subset U} \sqrt{\frac{\omega_{S}}{\omega_{S^{\prime}}}}\right)}\left\|\nabla v_{j}\right\|_{j, \omega, U}
$$

A direct consequence of Theorem 3.9 is the following lower bound.
Theorem 3.11. Let the bilinear form $a(\cdot, \cdot)$ fulfill Assumption 1 and consider a sequence of uniformly refined triangulations $\mathcal{T}_{j}$ and the respective sequence of nested spaces $\mathcal{V}_{j}$ of linear finite elements. Consider the sets $U=U_{T}$ and $S, S^{\prime} \in \mathcal{T}_{j}$ defined in Theorem 3.9 for every $T \in \mathcal{T}_{j-1}$. Now define

$$
\begin{equation*}
C_{j, \omega}:=\max _{T \in \mathcal{T}_{j-1}} \max _{\substack{S^{\prime}, S \in \mathcal{I}_{j} \\ S^{\prime}, S \in U_{T}}} \frac{\omega_{S}}{\omega_{S^{\prime}}} \quad \text { and } \quad C_{J, \text { max }, \omega}:=\max _{0 \leq j \leq J} C_{j, \omega} \tag{3.29}
\end{equation*}
$$

and let $\omega$ be such that we can find a minimal refinement level $j_{0}<J$ so that

$$
\begin{equation*}
2^{j}>C \sqrt{C_{J, \max , \omega}} \tag{3.30}
\end{equation*}
$$

holds for all $j \geq j_{0}$ and Theorem 3.9 is applicable for all $j \geq j_{0}$. Then there holds the estimate

$$
\begin{equation*}
\sum_{j=j_{0}}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \leq 4\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} C C_{J, \max , \omega} a(u, u) \tag{3.31}
\end{equation*}
$$

for $u_{j}$ defined in (3.8).

Proof. Using Lemma 3.7, compare (3.14), we establish the estimate

$$
\begin{aligned}
\sum_{j=j_{0}}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} & \leq 4\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} \sum_{j=j_{0}}^{J} 2^{2 j}\left\|v_{j}\right\|_{j, \omega}^{2} \\
& \leq 4\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} C C_{J, \max , \omega} \sum_{j=j_{0}}^{J} a\left(v_{j}, v_{j}\right)
\end{aligned}
$$

after squaring (3.17) and summation with respect to $U_{T}$, i.e., over all $T \in \mathcal{T}_{j-1}$, where the constant $C_{0}$ depends on the shape of the triangles of the initial triangulation only. Then, with (3.7) the assertion (3.31) follows.

Note that this estimate employs the minimal conditions on the coefficient function $\omega$ in the sense, that assuming a fixed a refinement level $J$ we may find such a $j_{0}$ for a rather large set of functions $\omega$. However, from a more practical point of view it is rather important to ensure the existence of $j_{0}$ independent of $J$ involving only information obtainable from $\omega$ directly. Hence, let us now focus on the formulation of conditions on the coefficient function $\omega$ such that the results of Theorem 3.11 still hold, but for $j_{0}$ independent of $J$.

In essence, we are looking for coefficient functions $\omega$ such that

$$
\begin{equation*}
C_{j, \omega}=\max _{T \in \mathcal{T}_{j-1}} \max _{\substack{S^{\prime}, S \in \mathcal{I}_{j} \\ S^{\prime}, S \in U_{T}}} \frac{\omega_{S}}{\omega_{S^{\prime}}}<\operatorname{diam}^{2}\left(U_{T}\right) \tag{3.32}
\end{equation*}
$$

starting from a certain level $j_{0}$. Hence, if we consider a non-negative function $\omega$ such that for all $x$ and $y$ with $\|x-y\|<2^{j_{0}}$ and $\omega(x)>\omega(y)$ the estimate

$$
\omega(x)<2^{2 j_{0}} w(y)
$$

holds, then condition (3.32) is fulfilled for all $j \geq j_{0}$. A trivial sufficient condition is of course

$$
C_{j, \omega} \leq C_{\omega}<\infty \quad \text { for all } j
$$

Also, it is obvious that continuous weights $\omega$ are admissible. Another simple condition which suffices is

$$
E_{j, \omega}:=\max _{T \in \mathcal{T}_{j}} \max _{x, y \in T}\left\|\frac{\omega(x)-\omega(y)}{\omega(y)}\right\| \rightarrow 0 \quad \text { for } j \rightarrow \infty
$$

since

$$
\frac{\omega_{S}}{\omega_{S^{\prime}}}=\frac{\int_{S} \omega}{\int_{S^{\prime}} \omega} \frac{\mu\left(S^{\prime}\right)}{\mu(S)}=\frac{\int_{S} \frac{\omega}{\omega(y)}}{\int_{S^{\prime}} \frac{\omega}{\omega(y)}} \frac{\mu\left(S^{\prime}\right)}{\mu(S)} \leq \frac{\left(1+E_{j, \omega}\right) \int_{S} 1}{\left(1-E_{j, \omega}\right) \int_{S^{\prime}} 1} \frac{\mu\left(S^{\prime}\right)}{\mu(S)}=\frac{\left(1+E_{j, \omega}\right)}{\left(1-E_{j, \omega}\right)}
$$

Furthermore, any piecewise constant coefficient function allows for the determination of $j_{0}$ independent of the number of jumps, their frequency and their location. Note that the estimate does not rely on any specific knowledge about the "resolution" of the jumps on a particular level or grid. Only the height of the maximal jump is relevant for $j_{0}$. Let $\epsilon^{-1}$ denote the height of the maximal jump, then choosing $j_{0}$ such that $\ln \left(\epsilon^{-1 / 2}\right)<j_{0}$ is sufficient.

Let us now consider $\omega \in A_{\infty}$, since we already required $\omega$ to be in the Muckenhoupt class in Lemma 2.2 to establish the norm equivalence between $\|\cdot\|_{\omega}$ and $\|\cdot\|_{j, \omega}$ in Theorem 2.4.

A function $\omega$ is in $A_{\infty}(\Omega)$ if and only if

$$
\frac{1}{\mu(B)} \int_{B} \omega d x \exp \left(\frac{1}{\mu(B)} \int_{B} \ln \left(\frac{1}{\omega}\right) d x\right) \leq A<\infty
$$

holds for all balls $B \subset \Omega$. Another characterization is given by: A weight $\omega$ is in $A_{\infty}$ if and only if for all balls $B$ and all $F \subset B$ with $\mu(F) \geq \alpha \mu(B)$ there exists a constant $\beta \in(0,1)$ such that

$$
\begin{equation*}
\int_{F} \omega \geq \beta \int_{B} \omega \tag{3.33}
\end{equation*}
$$

holds. Using this property, we obtain that any such weight fulfills our limited growth condition automatically. Consider $S, S^{\prime} \subset U$, since we employ uniformly refined triangulations, we have $\mu\left(S^{\prime}\right) \approx \mu(S) \geq \mu(U) / 16$. Now, let $B$ denote the smallest ball which completely contains $U$, i.e. $U \subset B$, then the estimate

$$
\mu\left(S^{\prime}\right) \approx \mu(S) \geq C_{\infty} \mu(B)
$$

with $C_{\infty}$ depends on the initial triangulation $\mathcal{T}_{0}$ only. Hence, with the characterization (3.33) given above we obtain

$$
\int_{S^{\prime}} \omega \geq \beta \int_{B} \omega
$$

and therefore

$$
\begin{equation*}
\frac{\omega_{S}}{\omega_{S^{\prime}}} \leq \frac{\int_{B} \omega}{\beta \int_{B} \omega} \frac{\mu\left(S^{\prime}\right)}{\mu(S)} \leq \frac{C_{2}}{\beta}:=\mathcal{C}_{\infty, \omega} \tag{3.34}
\end{equation*}
$$

Therefore, all quotients $\frac{\omega_{S}}{\omega_{S^{\prime}}}$ are bounded independent of the level $j$ by $\mathcal{C}_{\infty, \omega}$ for $\omega \in A_{\infty}$ and the results of Theorem 3.9 hold.

Remark 3.12. Note that there is a close connection between the Muckenhoupt class $A_{\infty}$ with the space BMO via the implications

$$
\omega \in A_{p} \subset A_{\infty} \quad \Rightarrow \quad \ln (\omega) \in \mathrm{BMO}
$$

and

$$
f \in \mathrm{BMO} \quad \Rightarrow \quad f=c \ln \omega \quad \text { with } \omega \in A_{\infty}
$$

Let us summarize our results in the following theorem.
Theorem 3.13. Let the bilinear form $a(\cdot, \cdot)$ fulfill Assumption 1 and consider a sequence of uniformly refined triangulations $\mathcal{T}_{j}$ and the respective sequence of nested spaces $\mathcal{V}_{j}$ of linear finite elements. Furthermore, let us consider $\omega \in A_{\infty}$. Then $C_{j, \omega}<\mathcal{C}_{\infty, \omega}$ holds for all $j$. Hence, there exists a minimal refinement level $j_{0}$ independent of $J$ such that the limited growth condition

$$
\begin{equation*}
2^{j}>C \sqrt{\mathcal{C}_{\infty, \omega}} \tag{3.35}
\end{equation*}
$$

is fulfilled for all $j \geq j_{0}$. Therefore, the equivalencies

$$
\begin{equation*}
a(u, u) \leq \sum_{j=0}^{j_{0}-1} a\left(v_{j}, v_{j}\right)+C_{3} \sum_{j=j_{0}}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2} \leq C_{4} a(u, u) \tag{3.36}
\end{equation*}
$$

and

$$
\begin{equation*}
a(u, u) \leq \sum_{j=0}^{j_{0}-1} a\left(v_{j}, v_{j}\right)+C_{5} \sum_{j=j_{0}}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \leq C_{6} a(u, u) \tag{3.37}
\end{equation*}
$$

hold with the constants

$$
\begin{equation*}
C_{3}:=\frac{32}{3} C_{2}, \quad C_{4}:=C_{3} 4\left(1+2 \gamma^{-1} C_{0}\right)^{2}\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} C \mathcal{C}_{\infty, \omega} \tag{3.38}
\end{equation*}
$$

and

$$
\begin{equation*}
C_{5}:=C_{3}\left(1+s \gamma^{-1} C_{0}\right)^{2}, \quad C_{6}:=C_{5} C \mathcal{C}_{\infty, \omega} \tag{3.39}
\end{equation*}
$$

where $C_{0}, C_{1}$, and $C_{2}$ depend on the initial triangulation $\mathcal{T}_{0}$ only and $\gamma$ is a constant determined by the $A_{\infty}$ property of the coefficient function $\omega$.

Let us now consider the remainder terms for the coarser levels $j=0, \ldots, j_{0}-$ 1. Since $j_{0}$ is independent of $J$ it is sufficient to establish an equivalence which explicitly involves $j_{0}$ (compare Theorem 3.2). Hence, we can employ a similar approach as in [23]. To this end, let us introduce the averaging operators $M_{j}$ similar to [23]. However, our averaging operators $M_{j}$ are based on the weighted scalar product $\langle\cdot, \cdot\rangle_{\omega}$, namely we use

$$
\begin{equation*}
M_{j}^{\omega} u:=\sum_{i=1}^{n_{j}} \frac{\left\langle u, \psi_{i}^{j}\right\rangle_{\omega}}{\left\langle 1, \psi_{i}^{j}\right\rangle_{\omega}} \psi_{i}^{j} \tag{3.40}
\end{equation*}
$$

where $\psi_{i}^{j}$ denotes the piecewise linear finite element basis function at vertex $i$ on level $j$ and $n_{j}=\# \mathcal{N}_{j}$ denotes the number of vertices in the triangulation $\mathcal{T}_{j}$. Similarly as in the proof of Lemma 3.7 we obtain

$$
\begin{equation*}
\left\|u_{j}\right\|_{\omega} \leq\left\|u-M_{j-1}^{\omega} v\right\|_{\omega} . \tag{3.41}
\end{equation*}
$$

Omitting the level superscript $j$ for the basis functions, we have on each triangle $T \in \mathcal{T}_{j}$ the equivalence

$$
\left.M_{j}^{\omega} u\right|_{T}=\sum_{\nu=1}^{3} \frac{\left\langle u, \psi_{i(T, \nu)}\right\rangle_{\omega}}{\left\langle 1, \psi_{i(T, \nu)}\right\rangle_{\omega}} \psi_{i(T, \nu)}
$$

where $i(T, \nu)$ indicates the respective vertex of $T$. For the ease of notation we use $\psi_{\nu}$ for $\psi_{i(T, \nu)}$ in the following. With this local formulation we directly obtain the estimate

$$
\left\|M_{j}^{\omega} u\right\|_{\omega, T}^{2} \leq\left(\sum_{\nu=1}^{3} \frac{\left|\left\langle u, \psi_{\nu}\right\rangle_{\omega}\right|}{\left\langle 1, \psi_{\nu}\right\rangle_{\omega}}\left\|\psi_{\nu}\right\|_{\omega, T}\right)^{2} \leq 3 \sum_{\nu=1}^{3}\left|\left\langle u, \psi_{\nu}\right\rangle_{\omega}\right|^{2} \frac{\left\|\psi_{\nu}\right\|_{\omega, T}^{2}}{\left\langle 1, \psi_{\nu}\right\rangle_{\omega}^{2}}
$$

With $(\psi(x))^{2} \leq \psi(x)$ for all basis functions $\psi$ we obtain

$$
\frac{\left\|\psi_{\nu}\right\|_{\omega, T}^{2}}{\left\langle 1, \psi_{\nu}\right\rangle_{\omega}^{2}}=\frac{\int_{T} \omega(x)\left|\psi_{\nu}(x)\right|^{2} d x}{\left(\int_{\operatorname{supp}\left(\psi_{\nu}\right)} \omega(x)\left|\psi_{\nu}(x)\right| d x\right)^{2}} \leq\left(\int_{\operatorname{supp}\left(\psi_{\nu}\right)} \omega(x)\left|\psi_{\nu}(x)\right| d x\right)^{-1}
$$

Hence, we have the local boundedness of the averaging operators $M_{j}^{\omega}$ with the estimate

$$
\begin{aligned}
\left|\left\langle u, \psi_{\nu}\right\rangle_{\omega}\right|^{2} & =\left|\int_{\operatorname{supp}\left(\psi_{\nu}\right)} u(x) \omega(x) \psi_{\nu}(x) d x\right|^{2} \\
& =\left|\int_{\operatorname{supp}\left(\psi_{\nu}\right)}\left(u(x) \sqrt{\omega(x) \psi_{\nu}(x)}\right) \sqrt{\omega(x) \psi_{\nu}(x)} d x\right|^{2} \\
& \leq \int_{\operatorname{supp}\left(\psi_{\nu}\right)} \omega(x) \psi_{\nu}(x) d x \int_{\operatorname{supp}\left(\psi_{\nu}\right)}|u(x)|^{2} \omega(x) \psi_{\nu}(x) d x .
\end{aligned}
$$

Lemma 3.14. For any $T \in \mathcal{T}_{j}$ let $U=\bigcup_{\nu=1}^{3} \operatorname{supp}\left(\psi_{\nu}\right)$ be the union of the supports of the three nodal basis functions $\psi_{\nu}=\psi_{i(T, \nu)}$ with $T \subset \operatorname{supp}\left(\psi_{\nu}\right)$. Then the estimate

$$
\left\|M_{j}^{\omega} u\right\|_{\omega, T}^{2} \leq 3 \sum_{\nu=1}^{3} \int_{\operatorname{supp}\left(\psi_{\nu}\right)} \omega(x)|u(x)|^{2} \psi_{\nu}(x) d x \leq 3 \int_{U} \omega(x)|u(x)|^{2} d x
$$

holds for any $u \in \mathcal{V}$ so that

$$
\begin{equation*}
\left\|M_{j}^{\omega} u\right\|_{\omega, T}^{2} \leq 3\|u\|_{\omega, U}^{2} \tag{3.42}
\end{equation*}
$$

Since $M_{j}^{\omega}$ preserves constants we obtain the estimate

$$
\begin{equation*}
\left\|M_{j}^{\omega} u-u\right\|_{\omega, T} \leq\left\|M_{j}^{\omega}\left(u-a_{T}\right)\right\|_{\omega, T}+\left\|u-a_{T}\right\|_{\omega, T} \leq(1+\sqrt{3})\left\|\left(u-a_{T}\right)\right\|_{\omega, U} \tag{3.43}
\end{equation*}
$$

with the help of this lemma. Now we use the Sobolev inequality in $\mathbb{R}^{2}$, i.e., the Taylor-remainder formula in integral form, which reads

$$
u(x)=\int_{U} \varphi(y) u(y) d y+R^{1} u(x)=\int_{U} \varphi(y) u(y) d y+\sum_{|\alpha|=1} \int_{U} K_{\alpha}(x, y) u^{(\alpha)}(y) d y
$$

where $\varphi$ is a smooth function with support in a ball $B \subset U$ and $\int \varphi(x) d x=1$, and the kernels $K_{\alpha}$ are given by

$$
K_{\alpha}(x, y)=(x-y)^{\alpha} \int_{0}^{1} s^{-3} \varphi\left(x+s^{-1}(y-x)\right) d s
$$

One can show that the estimate

$$
\left|\int_{U} K_{\alpha}(x, y) u^{(\alpha)}(y) d y\right| \leq C \frac{(\operatorname{diam}(U))^{2}}{\delta^{2}} \int_{U}\left|u^{(\alpha)}(z)\right|(x-z)^{\alpha}| | x-\left.z\right|^{-2} d z
$$

holds, with $C_{K}=C \frac{(\operatorname{diam}(U))^{2}}{\delta^{2}}$ depending only on the initial triangulation $\mathcal{T}_{0}$, where $C$ is an absolute constant independent of $u$, and $\delta$ is the radius of the ball $B$. Thus, choosing $a_{T}=\int_{U} \varphi(y) u(y) d y$ we obtain the estimate

$$
\left\|u(x)-a_{T}\right\|_{\omega, U}^{2} \leq C_{K}(\operatorname{diam}(U))^{2} \int_{U} \omega(x)\left|\sum_{|\alpha|=1} \int_{U}\right| u^{(\alpha)}(z) \| x-\left.\left.z\right|^{-2} d z\right|^{2} d x
$$

Now we introduce the kernel $G(x):=|x|^{-2}$ and rewrite the right-hand side as a singular integral with the notation $\left(u^{(\alpha)} *_{U} G\right)(x)=\int_{U} u^{(\alpha)}(z) G(x-z) d z$

$$
\left\|u(x)-a_{T}\right\|_{\omega, U}^{2} \leq C_{K}(\operatorname{diam}(U))^{2} \sum_{|\alpha|=1} \int_{U} \omega(x)\left|\left(u^{(\alpha)} *_{U} G\right)(x)\right|^{2} d x
$$

Then, using a result from harmonic analysis [18] for weights $\omega \in A_{2}$ (a sub-class of the Muckenhoupt class $A_{\infty}$ ) we obtain

$$
\int_{U} \omega(x)\left|\left(u^{(\alpha)} *_{U} G\right)(x)\right|^{2} d x \leq \mathcal{C}_{2, \omega} \int_{U} \omega(x)\left|u^{(\alpha)}\right|^{2} d x
$$

with $\mathcal{C}_{2, \omega}$ depending on $\omega$ only. Hence, with $\alpha=1$ and (3.43) we end up with

$$
\begin{equation*}
\left\|M_{j}^{\omega} u\right\|_{\omega, T} \leq(1+\sqrt{3})\left\|\left(u-a_{T}\right)\right\|_{\omega, U} \leq \mathcal{C}_{2, \omega} C_{K}(\operatorname{diam}(U))^{2}\|\nabla u\|_{\omega, U}^{2} \tag{3.44}
\end{equation*}
$$

Finally, we choose $v=M_{j}^{\omega} u$ in (3.41) and we obtain

$$
\left\|u_{j}\right\|_{j, \omega}^{2} \leq 4\left\|M_{j}^{\omega} u-u\right\|_{j, \omega}^{2} \leq \mathcal{C}_{2, \omega} 2^{-2 j} a(u, u)
$$

where we have used $\|\nabla u\|_{\omega}^{2}=a(u, u)$ after summation with respect to $U$, i.e. $T \in \mathcal{T}_{j}$.

Recall that we can determine $j_{0}$ for $\omega \in A_{\infty}$ due to the limited growth condition

$$
2^{j}>C \sqrt{\mathcal{C}_{\infty, \omega}}
$$

for all $j \geq j_{0}$ and we obtain

$$
j_{0}:=\ln \left(C \sqrt{\mathcal{C}_{\infty, \omega}}\right)
$$

Altogether, we can now establish our robust and optimal norm equivalencies in the following theorem.

Theorem 3.15. Let the bilinear form a $(\cdot, \cdot)$ fulfill Assumption 1 and consider a sequence of uniformly refined triangulations $\mathcal{T}_{j}$ and the respective sequence of nested spaces $\mathcal{V}_{j}$ of linear finite elements. Furthermore, let us consider $\omega \in A_{2} \subset A_{\infty}$. Then the equivalence

$$
\begin{equation*}
a(u, u) \leq C_{3} \sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{j, \omega}^{2} \leq C_{7} a(u, u) \tag{3.45}
\end{equation*}
$$

where

$$
C_{7}:=C_{3}\left(\ln \left(C \sqrt{\mathcal{C}_{\infty, \omega}}\right) \mathcal{C}_{2, \omega}+4\left(1+2 \gamma^{-1} C_{0}\right)^{2}\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2} C \mathcal{C}_{\infty, \omega}\right)
$$

$\gamma, \mathcal{C}_{2, \omega}$ and $\mathcal{C}_{\infty, \omega}$ depend on $\omega$ only, $C_{0}$ and $C_{3}$ depend on the initial triangulation $\mathcal{T}_{0}$ only. With respect to the $\|\cdot\|_{\omega}$ norm we have the equivalence

$$
\begin{equation*}
a(u, u) \leq C_{3}\left(1+2 \gamma^{-1} C_{0}\right)^{2} \sum_{j=0}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} \leq C_{8} a(u, u) \tag{3.46}
\end{equation*}
$$

where

$$
C_{8}:=C_{3}\left(1+2 \gamma^{-1} C_{0}\right)^{2}\left(\ln \left(C \sqrt{\mathcal{C}_{\infty, \omega}}\right) \mathcal{C}_{2, \omega}\left(1+\sqrt{\frac{3}{2} C_{0}}\right)^{2}+C \mathcal{C}_{\infty, \omega}\right)
$$

## 4. Robust and Optimal Preconditioners

Let us now focus on the design of appropriate preconditioners based on the established norm equivalencies. Here, however, we can derive a simple preconditioner based on the $\|\cdot\|_{\omega}$ norm only. The following standard construction for BPX-type preconditioners $B[3,22]$ does not work for the level-dependent norms $\|\cdot\|_{j, \omega}$.

Let us define the preconditioner $B$ via

$$
\begin{equation*}
\langle u, B u\rangle_{\omega}=a\left(Q_{0}^{\omega} u, Q_{0}^{\omega} u\right)+\sum_{j=1}^{J} 2^{2 j}\left\|u_{j}\right\|_{\omega}^{2} . \tag{4.1}
\end{equation*}
$$

Using the orthogonality property $\left\langle u_{j}, u_{l}\right\rangle_{\omega}=\delta_{j, l}$ for $u_{j}=\left(Q_{j}^{\omega}-Q_{j-1}^{\omega}\right) u$ we obtain the operator formulation

$$
B=A_{0} Q_{0}^{\omega}+\sum_{j=1}^{J} 2^{2 j}\left(Q_{j}^{\omega}-Q_{j-1}^{\omega}\right)
$$

Similarly, we obtain the inverse operator

$$
B^{-1}=A_{0}^{-1} Q_{0}^{\omega}+\sum_{j=1}^{J} 2^{-2 j}\left(Q_{j}^{\omega}-Q_{j-1}^{\omega}\right)
$$

Note that we can rewrite the sum of the differences of the projections also as a simple sum of the $Q_{j}^{\omega}$, i.e.

$$
\begin{aligned}
\sum_{j=1}^{J} 2^{-2 j}\left(Q_{j}^{\omega}-Q_{j-1}^{\omega}\right) & =2^{-2 J} Q_{J}^{\omega}-2^{-2} Q_{0}^{\omega}+\sum_{j=1}^{J-1} 2^{-2 j} Q_{j}^{\omega}-\sum_{j=1}^{J-1} 2^{-2 j-2} Q_{j}^{\omega} \\
& =2^{-2 J} Q_{J}^{\omega}-2^{-2} Q_{0}^{\omega}+\sum_{j=1}^{J-1}\left(1-2^{-2}\right) 2^{-2 j} Q_{j}^{\omega}
\end{aligned}
$$

With this we obtain the alternative representation of the operator $B^{-1}$ as

$$
B^{-1}=\left(A_{0}^{-1}-2^{-2}\right) Q_{0}^{\omega}+\sum_{j=1}^{J-1}\left(1-2^{-2}\right) 2^{-2 j} Q_{j}^{\omega}+2^{-2 J} Q_{J}^{\omega}
$$

Therefore, it is clear that the equivalent operator $\tilde{B}^{-1}$

$$
\tilde{B}^{-1}=A_{0}^{-1} Q_{0}^{\omega}+\sum_{j=1}^{J} 2^{-2 j} Q_{j}^{\omega}
$$

can also be used as a preconditioner. However, the orthogonal projections $Q_{j}^{\omega}$ are rather expensive to compute. Instead we will employ the spectrally equivalent averaging operators (3.40)

$$
M_{j}^{\omega} u=\sum_{i=1}^{n_{j}} \frac{\left\langle u, \psi_{i}^{j}\right\rangle_{\omega}}{\left\langle 1, \psi_{i}^{j}\right\rangle_{\omega}} \psi_{i}^{j}
$$

where $n_{j}=\# \mathcal{N}_{j}$ denotes the number of vertices of the triangulation $\mathcal{T}_{j}$. The operators $M_{j}^{\omega}$ satisfy

$$
\begin{equation*}
\gamma_{0}\left\langle u, M_{j}^{\omega} u\right\rangle_{\omega} \leq\left\langle u, Q_{j}^{\omega} u\right\rangle_{\omega} \leq \gamma_{1}\left\langle u, M_{j}^{\omega} u\right\rangle_{\omega} \text { uniformly in } j \tag{4.2}
\end{equation*}
$$

Finally, we obtain the substantially less expensive preconditioner $\hat{B}$ with

$$
\hat{B}^{-1}=A_{0}^{-1} Q_{0}^{\omega}+\sum_{j=1}^{J} 2^{-2 j} \sum_{i=1}^{n_{j}} \frac{\left\langle\cdot, \psi_{i}^{j}\right\rangle_{\omega}}{\left\langle 1, \psi_{i}^{j}\right\rangle_{\omega}} \psi_{i}^{j}
$$

It remains to show the spectral equivalence (4.2) which we obtain from the following lemma.

Lemma 4.1. Let the weight $\omega$ but such that (2.13) is fulfilled. Then the piecewise linear finite element functions $\psi_{l}^{(j)}$ form a Riesz-basis for the spaces $\mathcal{V}_{j}$ in the weighted norm $\|\cdot\|_{\omega}$ arising from the weighted scalar product (2.11); i.e., there holds the estimate

$$
\begin{equation*}
A \sum_{l=1}^{n_{j}}\left\langle 1, \psi_{l}^{(j)}\right\rangle_{\omega}\left|\alpha_{l}\right|^{2} \leq\left\|\sum_{l=1}^{n_{j}} \alpha_{l} \psi_{l}^{(j)}\right\|_{\omega}^{2} \leq B \sum_{l=1}^{n_{j}}\left\langle 1, \psi_{l}^{(j)}\right\rangle_{\omega}\left|\alpha_{l}\right|^{2} \tag{4.3}
\end{equation*}
$$

with constants $A$ and $B$ independent of the level $j$.

Proof. The inequality on the right-hand side follows from

$$
\left\|\sum_{l=1}^{n_{j}} \alpha_{l} \psi_{l}^{(j)}\right\|_{\omega}^{2} \leq\left\|\left(\sum_{l=1}^{n_{j}}\left|\alpha_{l}\right|^{2} \psi_{l}^{(j)}\right)^{1 / 2}\right\|_{\omega}^{2}=\sum_{l \in \mathcal{N}_{j}}\left|\alpha_{l}\right|^{2}\left\langle 1, \psi_{l}^{(j)}\right\rangle_{\omega}
$$

With the help of Theorem 2.4 we obtain

$$
\begin{aligned}
\left\|\sum_{l=1}^{n_{j}} \alpha_{l} \psi_{l}^{(j)}\right\|_{\omega}^{2} & \geq\left(1+2 \gamma^{-1} C_{0}\right)\left\|\sum_{l=1}^{n_{j}} \alpha_{l} \psi_{l}^{(j)}\right\|_{j, \omega}^{2} \\
& =\left(1+2 \gamma^{-1} C_{0}\right) \sum_{T \in \mathcal{T}_{j}} \omega_{T} \int_{T}\left|\sum_{\nu=1}^{3} \alpha_{\nu, T} \psi_{\nu, T}\right|^{2}
\end{aligned}
$$

for the inequality on the left-hand side, where the $\psi_{\nu, T}$ denote the three nodal functions of level $j$ with $\operatorname{supp}\left(\psi_{\nu, T}\right) \supset T$. By the well-known stability properties of the $\psi_{l}^{(j)}$ there holds (cf. [23])

$$
\int_{T}\left|\sum_{\nu=1}^{3} \alpha_{\nu, T} \psi_{\nu, T}^{(j)}\right|^{2} \geq \sum_{\nu=1}^{3}\left|\alpha_{\nu, T}\right|^{2}\left\langle 1, \psi_{\nu, T}\right\rangle_{T}
$$

Since the support of each $\psi_{i}^{(j)}$ consists only of a finite number of triangles it follows further that

$$
\begin{align*}
\left\|\sum_{l=1}^{n_{j}} \alpha_{l} \psi_{l}^{(j)}\right\|_{\omega}^{2} & \geq\left(1+2 \gamma^{-1} C_{0}\right) \sum_{T \in \mathcal{T}_{j}} \omega_{T} \sum_{\nu=1}^{3}\left|\alpha_{\nu, T}\right|^{2}\left\langle 1, \psi_{\nu, T}\right\rangle_{T}  \tag{4.4}\\
& \geq C\left(1+2 \gamma^{-1} C_{0}\right) \sum_{l=1}^{n_{j}}\left|\alpha_{l}\right|^{2}\left\langle 1, \psi_{l}^{(j)}\right\rangle_{j, \omega}
\end{align*}
$$

with an absolute constant $C$. Now, with the $L_{1}$-counterpart of Theorem 2.4 we obtain the estimate

$$
\left\|\sum_{l=1}^{n_{j}} \alpha_{l} \psi_{l}^{(j)}\right\|_{\omega}^{2} \geq \bar{C} \sum_{l=1}^{n_{j}}\left|\alpha_{l}\right|^{2}\left\langle 1, \psi_{l}^{(j)}\right\rangle_{\omega}
$$

Note that this equivalence involves the same constants as Theorem 2.4, i.e., involves some information on the coefficient function $\omega$ and holds for $\omega \in A_{\infty}$. With the help of this lemma we obtain (4.2) by the following consideration, see [23]. With $Q_{j}^{\omega} u=\alpha_{l} \psi_{l}^{(j)}$ we have the equivalence

$$
\left\langle u, Q_{j}^{\omega} u\right\rangle_{\omega}=\left\langle Q_{j}^{\omega} u, Q_{j}^{\omega} u\right\rangle_{\omega}=\left\langle\vec{\alpha}, G_{j} \vec{\alpha}\right\rangle_{\mathbb{R}^{n_{j}}}
$$

where $G_{j}:=\left(\left\langle\psi_{l}^{(j)}, \psi_{k}^{(j)}\right\rangle_{\omega}\right)_{l, k}$ denotes the Gram matrix. Now observe that with $D_{j}:=\operatorname{diag}\left(\left\langle 1, \psi_{l}^{(j)}\right\rangle_{\omega}\right)$ and $\hat{u}:=Q_{j}^{\omega} u$, we also have $\left\langle G_{j} \vec{\alpha}, D_{j}^{-1} G_{j} \vec{\alpha}\right\rangle_{\mathbb{R}^{n_{j}}}=\left\langle\hat{u}, M_{j}^{\omega} \hat{u}\right\rangle_{\omega}=\left\langle Q_{j}^{\omega} u, M_{j}^{\omega} Q_{j}^{\omega} u\right\rangle_{\omega}=\left\langle Q_{j}^{\omega} u, M_{j}^{\omega} u\right\rangle_{\omega}=\left\langle u, M_{j}^{\omega} u\right\rangle_{\omega}$. since $Q_{j}^{\omega}$ is an orthogonal projection. Finally, with the notation $\vec{\beta}:=G_{j} \vec{\alpha}$ and Lemma 4.1 we obtain

$$
\left\langle\vec{\alpha}, G_{j} \vec{\alpha}\right\rangle_{\mathbb{R}^{n_{j}}}=\left\langle G_{j}^{-1} \vec{\beta}, \vec{\beta}\right\rangle_{\mathbb{R}^{n_{j}}} \asymp\left\langle\vec{\beta}, D_{j}^{-1} \vec{\beta}\right\rangle_{\mathbb{R}^{n_{j}}}=\left\langle G_{j} \vec{\alpha}, D_{j}^{-1} G_{j} \vec{\alpha}\right\rangle_{\mathbb{R}^{n_{j}}}
$$

which is only the discrete form of (4.2). This completes the construction of our preconditioner based on the $\|\cdot\|_{\omega}$ norm.

Remark 4.2. Note, however, that this preconditioner still involves the weights $\omega_{T}$ with $T \in \mathcal{T}_{J}$ for all levels $j$. Hence, the evaluation of the $\left\langle 1, \psi_{l}^{(j)}\right\rangle_{\omega}$ terms is not trivial.

The development of a preconditioner based on the level-dependent norms $\|\cdot\|_{j, \omega}$ which should give a superior performance, i.e., should lead to a smaller condition number, is an open problem. Already the definition of a simple linear operator $\mathcal{B}$ similar to (4.1) is not trivial since the terms on the right-hand side now involve the level-dependent norms $\|\cdot\|_{j, \omega}$ and the orthogonality property holds for the norm $\|\cdot\|_{\omega}$ on the finest level only. The derivation of such a preconditioner is subject of current research and will be addressed in a forthcoming paper.

## 5. Concluding Remarks

We presented two optimal and (almost) robust norm equivalencies based on certain weighted norms for diffusion problems $-\nabla \omega \nabla u=f$ in two space dimensions with a scalar diffusion coefficient $\omega$. We only require $\omega$ to be in the Muckenhoupt class $A_{2} \subset A_{\infty}$ to obtain our optimal bounds which involve some information on the boundesness of the local variation of the coefficient function $\omega$. This covers all piecewise constant functions independent of the location of jumps, their number or their frequency. This is in contrast to previous results, which require the resolution of the jumps on a particular level, i.e. the coarsest level.

Based on the weighted norm $\|\cdot\|_{\omega}$ which involves the evaluation of the piecewise constant approximation of the weights $\omega_{T}$ with respect to the finest triangulation $\mathcal{T}_{J}$ we have furthermore presented a simple and (nearly) robust BPX-type preconditioner. The design of a preconditioner based on the more sophisticated level-dependent norms $\|\cdot\|_{j, \omega}$ is more involved and subject of current research.

## References

1. R. E. Alcouffe, A. Brandt, J. E. Dendy and J. W. Painter, The Multi-Grid Method for the Diffusion Equation with Strongly Discontinuous Coefficients, SIAM J. Sci. Comput. 2 (1981), 430-454.
2. D. Braess, Finite Elements: Theory, Fast Solvers, and Applications in Solid Mechanics, Cambridge University Press, 2001.
3. J. H. Bramble, J. E. Pasicak, and J. Xu, Parallel Multilevel Preconditioners, Math. Comp. 55 (1990), no. 191, 1-22.
4. A. Brandt, Multi-Level Adaptive Technique (MLAT) for Fast Numerical Solution to Boundary Value Problems, Proc. of the Third Int. Conf. on Numerical Methods in Fluid Mechanics, Univ. Paris 1972 (New York, Berlin, Heidelberg) (H. Cabannes and R. Teman, eds.), Springer, 1973.
5. $\qquad$ , Multi-Level Adaptive Solutions to Boundary-Value Problems, Math. Comp. 31 (1977), 333-390.
6. $\qquad$ , Algebraic Multigrid Theory: The Symmetric Case, Preliminary Proceedings for the International Multigrid Conference (Copper Mountain, Colorado), April 1983.
7. __ Algebraic Multigrid Theory: The Symmetric Case, Appl. Math. Comput. 19 (1986), 23-56.
8. A. Brandt, S. F. McCormick, and J. W. Ruge, Algebraic Multigrid for Automatic Multigrid Solutions with Application to Geodetic Computations, Technical Report, Institute for Computational Studies, Fort Collins, Colorado, October 1982.
9. , Algebraic Multigrid for Sparse Matrix Equations, Sparsity and Its Applications (D. J. Evans, ed.), Cambridge University Press, 1984.
10. J. E. Dendy, Black Box Multigrid, J. Comput. Phys. 48 (1982), 366-386.
11. M. Dryja, M. V. Sarkis, and O. B. Widlund, Multilevel Schwarz Methods for Elliptic Problems with Discontinuous Cefficients in Three Dimensions, Numer. Math. 72 (1996), 313-348.
12. T. Grauschopf, M. Griebel, and H. Regler, Additive Multilevel-Preconditioners based on Bilinear Interpolation, Matrix Dependent Geometric Coarsening and Algebraic Multigrid Coarsening for Second Order Elliptic PDEs, Applied Numerical Mathematics 23(1) (1997), 63-96.
13. W. Hackbusch, Ein iteratives Verfahren zur schnellen Auflösung elliptischer Randwertprobleme, Tech. Report 76-12, Mathematisches Institut, Universität zu Köln, 1976.
14. __ A Fast Numerical Method for Elliptic Boundary Value Problems with Variable Coefficients, 2nd GAMM-Conf. Numer. Meth. Fl. Mech. (Köln) (E. H. Hirschel and W. Geller, eds.), Deutsche Forschungs- und Versuchsanstalt für Luft- und Raumfahrt, 1977, pp. 50-57.
15. P. Oswald, On the Robustness of the BPX-Preconditioner with Respect to Jumps in the Coefficients, Math. Comp. 68 (1999), 633-650.
16. J. W. Ruge and K. Stüben, Efficient Solution of Finite Difference and Finite Element Equations by Algebraic Multigrid, Multigrid Methods for Integral and Differential Equations (D. J. Paddon and H. Holstein, eds.), The Institute of Mathematics and its Applications Conference Series, Clarendon Press, 1985.
17. K. Scherer, Weighted Jackson and Bernstein Inequalities, Preprint 120, Sonderforschungsbereich 611, 2002, submitted to Constructive Approximation.
18. E. M. Stein, Harmonic Analysis: Real-Variable Methods, Orthogonality and Oscillatory Integrals, Princeton University Press, 1993.
19. U. Trottenberg, C. W. Osterlee, and A. Schüller, Multigrid, Appendix A: An Introduction to Algebraic Multigrid by K. Stüben, pp. 413-532, Academic Press, San Diego, 2001.
20. W. L. Wan, T. F. Chan, and B. Smith, An Energy-Minimizing Interpolation for Robust Multigrid Methods, SIAM J. Sci. Comput. 21 (2000), no. 4, 1632-1649.
21. J. Wloka, Partielle Differentialgleichungen, Teubner, 1982.
22. J. Xu, Iterative Methods by Space Decomposition and Subspace Correction, SIAM Review 34 (1992), no. 4, 581-613.
23. H. Yserentant, Two Preconditioners Based on the Multi-level Splitting of Finite Element Spaces, Numer. Math. 58 (1990), 163-184.
24. P. M. Zeeuw, Matrix-Dependent Prolongations and Restrictions in a Black-Box Multigrid Solver, J. Comput. Appl. Math. 33 (1990), 1-27.
25. , Acceleration of Iterative Methods by Coarse Grid Corrections, Ph.D. thesis, University of Amsterdam, 1997.

Institut für Numerische Simulation, Universität Bonn, Germany
E-mail address: griebel@ins.uni-bonn.de
Institut für Angewandte Mathematik, Universität Bonn, Germany
E-mail address: scherer@iam.uni-bonn.de
Institut für Numerische Simulation, Universität Bonn, Germany
E-mail address: m.a.schweitzer@ins.uni-bonn.de


[^0]:    1991 Mathematics Subject Classification. Primary 65N55, 65F35; Secondary 65N30, 65F10.
    Key words and phrases. Norm equivalency, multilevel method, preconditioning, robustness.
    The authors were supported in part by the Sonderforschungsbereich 611 Singuläre Phänomene und Skalierung in Mathematischen Modellen sponsored by the Deutsche Forschungsgemeinschaft.

